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The L^2 Aeppli-Bott-Chern Hilbert complex [☆]



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ABSTRACT

We analyse the L^2 Hilbert complexes naturally associated to a non-compact complex manifold, namely the ones which originate from the Dolbeault and the Aeppli-Bott-Chern complexes. In particular we define the L^2 Aeppli-Bott-Chern Hilbert complex and examine its main properties on general Hermitian manifolds, on complete Kähler manifolds and on Galois coverings of compact complex manifolds. The main results are achieved through the study of self-adjoint extensions of various differential operators whose kernels, on compact Hermitian manifolds, are isomorphic to either Aeppli or Bott-Chern cohomology.

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1. Introduction

A *Hilbert complex* consists of a complex of mutually orthogonal Hilbert spaces \mathcal{H}_i along with linear operators $D_i : \mathcal{H}_i \rightarrow \mathcal{H}_{i+1}$ which are *densely defined* (D_i is defined on a domain $\mathcal{D}(D_i)$ which is a dense subspace of \mathcal{H}_i) and *closed* (the graph of D_i is closed in $\mathcal{H}_i \times \mathcal{H}_{i+1}$) of the form

$$0 \longrightarrow \mathcal{H}_0 \xrightarrow{D_0} \mathcal{H}_1 \longrightarrow \cdots \longrightarrow \mathcal{H}_{n-1} \xrightarrow{D_{n-1}} \mathcal{H}_n \longrightarrow 0$$

such that $\text{im } D_i \subseteq \ker D_{i+1}$. The notion of *Hilbert complexes* was systematically analysed by Brüning and Lesch in [11]. In this paper we are interested in studying the Hilbert complexes naturally associated to a non-compact complex manifold without boundary.

Recall that on a complex manifold M the exterior derivative on (p, q) -forms splits as

$$d = \partial + \bar{\partial} : A^{p,q} \longrightarrow A^{p+1,q} \oplus A^{p,q+1},$$

where

$$\partial : A^{p,q} \longrightarrow A^{p+1,q}, \quad \bar{\partial} : A^{p,q} \longrightarrow A^{p,q+1},$$

thus the relation $d^2 = 0$ implies $\partial^2 = \bar{\partial}^2 = \partial\bar{\partial} + \bar{\partial}\partial = 0$. Therefore the following cohomologies

$$H_{\bar{\partial}}^{\bullet,\bullet} = \frac{\ker \bar{\partial}}{\text{im } \bar{\partial}}, \quad H_{BC}^{\bullet,\bullet} = \frac{\ker \partial \cap \ker \bar{\partial}}{\text{im } \partial \bar{\partial}}, \quad H_A^{\bullet,\bullet} = \frac{\ker \partial \bar{\partial}}{\text{im } \partial + \text{im } \bar{\partial}},$$

called respectively *Dolbeault*, *Bott-Chern* and *Aeppli cohomology spaces* are well-defined. The elliptic complexes to which these cohomologies are associated are respectively the *Dolbeault complex*

$$\dots \longrightarrow A^{p,q-1} \xrightarrow{\bar{\partial}} A^{p,q} \xrightarrow{\bar{\partial}} A^{p,q+1} \longrightarrow \dots$$

and the *Aeppli-Bott-Chern complex*, or *ABC complex* for short,

$$\dots \longrightarrow A^{p-1,q-2} \oplus A^{p-2,q-1} \xrightarrow{\bar{\partial} \oplus \partial} A^{p-1,q-1} \xrightarrow{\partial \bar{\partial}} A^{p,q} \xrightarrow{\partial + \bar{\partial}} A^{p+1,q} \oplus A^{p,q+1} \longrightarrow \dots$$

where $\bar{\partial} \oplus \partial$ operates on $A^{p-1,q-2} \oplus A^{p-2,q-1}$ as $\bar{\partial}$ on $A^{p-1,q-2}$ plus ∂ on $A^{p-2,q-1}$. In particular the *Aeppli complex* is given by the differentials $\bar{\partial} \oplus \partial$ and $\partial \bar{\partial}$, while the *Bott-Chern complex* is given by the differentials $\partial \bar{\partial}$ and $\partial + \bar{\partial}$. The ABC complex first appeared in [10].

If we fix a Hermitian metric g on the complex manifold M with complex dimension n , then ∂ and $\bar{\partial}$ have L^2 formal adjoint operators

$$\partial^* : A^{p,q} \longrightarrow A^{p-1,q}, \quad \bar{\partial}^* : A^{p,q} \longrightarrow A^{p,q-1},$$

defined by $\partial^* := - * \bar{\partial}^*$ and $\bar{\partial}^* := - * \partial^*$, where $* : A^{p,q} \rightarrow A^{n-q,n-p}$ is the \mathbb{C} -linear Hodge operator with respect to g . The Laplacian which is naturally associated to the Dolbeault complex is the *Dolbeault Laplacian*

$$\Delta_{\bar{\partial}} = \bar{\partial} \bar{\partial}^* + \bar{\partial}^* \bar{\partial},$$

which is elliptic and formally self-adjoint, while to the Bott-Chern complex (for the Aeppli complex the situation is similar) there are multiple naturally associated Laplacians. The first is obtained as the standard Laplacian which is associated to any complex, namely

$$\Delta_{BC} = \partial \bar{\partial} \bar{\partial}^* \partial^* + \partial^* \partial + \bar{\partial}^* \bar{\partial},$$

which is formally self-adjoint but not elliptic [39]. Kodaira and Spencer, in the proof of the stability of the Kähler condition under small deformations of the complex structure [26], introduced another differential operator

$$\tilde{\Delta}_{BC} = \partial \bar{\partial} \bar{\partial}^* \partial^* + \bar{\partial}^* \partial^* \partial \bar{\partial} + \partial^* \bar{\partial} \bar{\partial}^* \partial + \bar{\partial}^* \partial \partial^* \bar{\partial} + \partial^* \partial + \bar{\partial}^* \bar{\partial},$$

usually referred as the *Bott-Chern Laplacian*, which is elliptic and formally self-adjoint and, when the manifold is compact, has the same kernel as Δ_{BC} . We will be also interested in a third Laplacian which is associated to the Bott-Chern complex, that is

$$\square_{BC} = \partial \bar{\partial} \bar{\partial}^* \partial^* + (\partial^* \partial + \bar{\partial}^* \bar{\partial})^2,$$

which is the elliptic and formally self-adjoint operator naturally associated to the Bott-Chern complex viewed as an elliptic complex [43]. When the manifold is compact \square_{BC} has the same kernel of Δ_{BC} and $\tilde{\Delta}_{BC}$.

In order to obtain a Hilbert complex starting from a geometric complex like the ones just mentioned, one first restricts every differential P of the complex to an operator P_0 defined on the space of smooth compactly supported forms, and then extends the restricted differential P_0 to a densely defined and closed operator on the space of L^2 forms with respect to a chosen metric: the two most important closed extensions are called *strong* and *weak*, denoted respectively by P_s and P_w , which correspond respectively to the minimal and the maximal closed extensions defined in a distributional sense of the operator P_0 . The L^2 Hodge theory related to the Hilbert complex originated from the Dolbeault complex, which from now on will be called L^2 Dolbeault Hilbert complex, has been studied even before [11], in parallel with the development of the L^2 Hodge theory associated to the de Rham complex: we refer to [2,24,19,12,14,16,21] as a partial list of milestones on this subject.

On the other hand, the study of a Hilbert complex arising from the ABC complex is absent in the literature. The aim of this paper is to define an L^2 Aeppli-Bott-Chern Hilbert complex, L^2 ABC Hilbert complex for short, and to establish its fundamental properties. In this setting the main difficulties arise from the fact that in the ABC complex the differential $\partial\bar{\partial}$ is of second order, while in the classical Hodge-de Rham or Dolbeault complexes all the differentials are of first order. For example, it is well known that on a complete Riemannian or Hermitian manifold, $P_s = P_w$ and $P_s^* = P_w^*$ (where P^* is the formal adjoint of P) for many interesting first order differential operators like $P = d, \partial, \bar{\partial}$ [2, Proposition 5]. Furthermore, the integer powers of $P + P^*$ (including the Dolbeault Laplacian $\Delta_{\bar{\partial}}$ and the Hodge Laplacian) are essentially self-adjoint [12, Section 3.B], *i.e.*, they have a unique self-adjoint extension.

The main results of this paper are contained in Sections 7, 8, 9. In Section 7 we define the L^2 ABC Hilbert complex

$$\begin{array}{c}
 L^2\Lambda^{p-1,q-2} \oplus L^2\Lambda^{p-2,q-1} \\
 \downarrow (\bar{\partial} \oplus \partial)_a \\
 L^2\Lambda^{p-1,q-1} \\
 \downarrow \partial\bar{\partial}_b \\
 L^2\Lambda^{p,q} \\
 \downarrow (\partial + \bar{\partial})_c \\
 L^2\Lambda^{p+1,q} \oplus L^2\Lambda^{p,q+1}
 \end{array}$$

where $a, b, c \in \{s, w\}$ with $a \leq b \leq c$ denote either strong or weak extensions with the order relation $s \leq w$. The associated L^2 Aeppli and Bott-Chern cohomologies are

$$L^2 H_{A,ab}^{p-1,q-1} := \frac{L^2 \Lambda^{p-1,q-1} \cap \ker \partial \bar{\partial}_b}{L^2 \Lambda^{p-1,q-1} \cap \text{im}(\bar{\partial} \oplus \partial)_a}, \quad L^2 H_{BC,bc}^{p,q} := \frac{L^2 \Lambda^{p,q} \cap \ker(\partial + \bar{\partial})_c}{L^2 \Lambda^{p,q} \cap \text{im} \partial \bar{\partial}_b},$$

while the reduced L^2 Aeppli and Bott-Chern cohomologies are defined as

$$L^2 \bar{H}_{A,ab}^{p-1,q-1} := \frac{L^2 \Lambda^{p-1,q-1} \cap \ker \partial \bar{\partial}_b}{L^2 \Lambda^{p-1,q-1} \cap \overline{\text{im}(\partial \oplus \bar{\partial})_a}}, \quad L^2 \bar{H}_{BC,bc}^{p,q} := \frac{L^2 \Lambda^{p,q} \cap \ker(\partial + \bar{\partial})_c}{L^2 \Lambda^{p,q} \cap \overline{\text{im} \partial \bar{\partial}_b}}.$$

We now list the main results of the paper (we state them just for the Bott-Chern case since the Aeppli case is analogous):

Theorem 1.1 (Section 7). *Given a Hermitian manifold (M, g) , we set the space of L^2 Bott-Chern harmonic forms to be $L^2 \mathcal{H}_{BC,bc}^{p,q} := \ker(\partial + \bar{\partial})_c \cap \ker \bar{\partial}^* \partial_b^*$, where $b' = s$ if $b = w$ and $b' = w$ if $b = s$.*

- The space of L^2 (p, q) -forms has an orthogonal decomposition

$$L^2 \Lambda^{p,q} = L^2 \mathcal{H}_{BC,bc}^{p,q} \oplus \overline{\text{im} \partial \bar{\partial}_b} \oplus \overline{\text{im}(\partial^* \oplus \bar{\partial}^*)_{c'}}.$$

- The reduced L^2 Bott-Chern cohomology is isomorphic to the space of L^2 Bott-Chern harmonic forms, i.e., $L^2 \bar{H}_{BC,bc}^{p,q} \simeq L^2 \mathcal{H}_{BC,bc}^{p,q}$.
- The L^2 Bott-Chern cohomology $L^2 H_{BC,bc}^{p,q}$ can be computed via a natural smooth sub-complex.
- If \square_{BC} is essentially self adjoint, then $\partial \bar{\partial}_s = \partial \bar{\partial}_w$.
- There exists a diagram of maps between spaces of reduced L^2 Bott-Chern, Dolbeault, de Rham, ∂ and Aeppli cohomologies.

The first three points are an application of the general theory of [11] to the specific case of the L^2 ABC complex, the fourth point is a generalisation of [11, Lemma 3.8], while the last point generalises the well known diagram of maps between usual cohomology spaces [13].

In Section 8, we focus on complete Kähler manifolds. Recall that a Hermitian metric is called Kähler when its fundamental form is closed, and that on a compact complex manifold endowed with a Kähler metric the kernels of the Bott-Chern, Aeppli and Dolbeault Laplacians coincide. By Hodge theory, this implies that Bott-Chern, Aeppli and Dolbeault cohomology spaces are isomorphic and the $\partial \bar{\partial}$ -Lemma holds. We prove the following

Theorem 1.2 (Section 8). *Let (M, g) be a complete Kähler manifold.*

- The space of L^2 Bott-Chern harmonic forms coincides with the space of L^2 Aeppli, Dolbeault, Hodge harmonic forms.
- All the maps in the diagram of Theorem 1.1 are isomorphisms.

- There exists an L^2 reduced $\partial\bar{\partial}$ -Lemma: if α is an L^2 k -form lying in $\ker \partial_w \cap \ker \bar{\partial}_w$, then

$$\begin{aligned} \alpha \in \overline{\text{im } \partial\bar{\partial}_s} &\iff \alpha \in \overline{\text{im } d_s} \iff \alpha \in \overline{\text{im } \partial_s} \\ &\iff \alpha \in \overline{\text{im } \bar{\partial}_s} \iff \alpha \in \overline{\text{im } \partial_s} + \overline{\text{im } \bar{\partial}_s}. \end{aligned}$$

- If the unique self-adjoint extension of $\Delta_{\bar{\partial}}$ has a spectral gap, then the reduced and unreduced L^2 Bott-Chern cohomologies coincide and $\partial\bar{\partial}_s = \partial\bar{\partial}_w$.

The first three points are generalisations of similar well-known results in the compact Kähler setting, while the last point is inspired by [21, Theorem 1.4.A]. The results are all obtained through the study of suitable self adjoint extensions of the Bott-Chern Laplacian $\tilde{\Delta}_{BC}$.

Finally, in Section 9, we study the special setting of a Galois covering of a compact complex manifold $\pi : \tilde{M} \rightarrow M \simeq \tilde{M}/\Gamma$. By [5, Proposition 3.1], any lift to the covering \tilde{M} of an elliptic and formally self-adjoint operator on the compact manifold M is essentially self-adjoint, therefore by the fourth point of Theorem 1.1 we obtain $\partial\bar{\partial}_s = \partial\bar{\partial}_w$ on \tilde{M} . This ultimately allows us to define L^2 Bott-Chern numbers $h_{BC,\Gamma}^{p,q}(M)$ and L^2 Aepli numbers $h_{A,\Gamma}^{p,q}(M)$ of the Galois covering (independently from the choice of the metric on the compact manifold), generalising the usual Bott-Chern and Aepli numbers of a compact complex manifold. E.g., $h_{BC,\Gamma}^{p,q}(M)$ is defined as the Von Neumann dimension of $L^2\mathcal{H}_{BC,sw}^{p,q}$ on \tilde{M} . Denoting by $h_{\partial,\Gamma}^{p,q}(M)$ and $h_{\bar{\partial},\Gamma}^{p,q}(M)$ the L^2 Hodge numbers, we prove the following inequality.

Theorem 1.3. *Given a Galois covering of a compact complex manifold $\pi : \tilde{M} \rightarrow M \simeq \tilde{M}/\Gamma$, it holds that*

$$h_{\partial,\Gamma}^{p,q}(M) + h_{\bar{\partial},\Gamma}^{p,q}(M) \leq h_{A,\Gamma}^{p,q}(M) + h_{BC,\Gamma}^{p,q}(M).$$

This is a generalisation of the same inequality for the usual Hodge, Bott-Chern and Aepli numbers on a compact complex manifold, established by Angella and Tomassini in [3, Theorem A].

We remark that all three Laplacians Δ_{BC} , $\tilde{\Delta}_{BC}$ and \square_{BC} enter in play in different arguments in our treatment. The operator Δ_{BC} is needed to apply the general theory of [11] to our case; on the other hand $\tilde{\Delta}_{BC}$ is useful on Kähler manifolds to prove that L^2 Bott-Chern harmonic forms coincide with L^2 Dolbeault harmonic forms in Theorem 1.2; finally \square_{BC} is fundamental when dealing with Galois coverings of compact complex manifolds: it is elliptic and formally self-adjoint (allowing us to apply [5, Proposition 3.1]) and it is naturally defined from the ABC complex (allowing us to prove the fourth point of Theorem 1.1).

We refer the reader to the following recent papers on L^2 Hodge and cohomology theory on non-compact (almost) Hermitian manifolds [7,8,23,38,35,25,29,30,32]. Furthermore,

we mention [33,34], where Tomassini and the second author prove a characterisation of L^2 smooth forms which are in the kernel of the Bott-Chern Laplacian on special families of Stein manifolds and of complete Hermitian manifolds; [42], where Tan, Wang and Zhou study L^2 Bott-Chern and Aeppli symplectic cohomologies and harmonic forms on complete non-compact almost Kähler manifolds with bounded geometry and introduce the L^2 dd^A Lemma; [9], where Bei, Diverio, Eyssidieux and Trapani generalise the notion of Kähler hyperbolicity introduced by Gromov in [21], proving a spectral gap result for the Dolbeault Laplacian under suitable modification.

The paper is structured in the following way. In Section 2 we define differential operators, Lebesgue spaces and elliptic complexes on manifolds. In Section 3 we describe the elliptic complexes naturally defined on a complex manifold and their associated Laplacians. In Section 4 we give a brief survey of the theory of unbounded linear operators on Hilbert spaces, including a statement of the Spectral Theorem and a study of the spectral gap condition for a positive self-adjoint operator. In Section 5 we define the minimal and maximal closed extensions of differential operators, collecting the main properties which will be used in the following sections. In Section 6 we recall the fundamental definitions concerning the L^2 Dolbeault Hilbert complex. Finally, in Section 7, 8, 9 we prove the main results of the paper. We conclude in Section 10 with some open questions and final remarks.

In this manuscript, our aim was to include a good amount of the necessary preliminaries, with the intention of making it readable for mathematicians who are not familiar with L^2 Hodge theory and spectral theory.

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2. Differential operators and elliptic complexes

In this section we follow mainly [15, Chapter VI, Section 1] and [43, Section I]. Let M be a differentiable manifold of dimension m , and let E_1, E_2 be \mathbb{C} -vector bundles over M , with rank $E_i = r_i$, $i = 1, 2$. Denote by $\Gamma(M, E_i)$ the spaces of smooth sections $M \rightarrow E_i$. A \mathbb{C} -linear *differential operator* of order l from E_1 to E_2 is a \mathbb{C} -linear operator $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ locally given by

$$Pu(x) = \sum_{\|\alpha\| \leq l} a_\alpha(x) D^\alpha u(x) \quad \forall u \in \Gamma(M, E_1),$$

where $\alpha = (\alpha_1, \dots, \alpha_m)$ is a multi-index with norm $\|\alpha\| = \alpha_1 + \dots + \alpha_m$, the functions

$$a_\alpha(x) = (a_{\alpha ij}(x))_{1 \leq i \leq r_2, 1 \leq j \leq r_1}$$

are $r_2 \times r_1$ matrices with smooth coefficients, $E_{1|\Omega} \simeq \Omega \times \mathbb{C}^{r_1}$, $E_{2|\Omega} \simeq \Omega \times \mathbb{C}^{r_2}$ are trivialized locally on some open chart $\Omega \subset M$ equipped with local coordinates x^1, \dots, x^m ,

$$D^\alpha = (\partial/\partial x^1)^{\alpha_1} \dots (\partial/\partial x^m)^{\alpha_m},$$

and $u = (u_j)_{1 \leq j \leq r_1}$, $D^\alpha u = (D^\alpha u_j)_{1 \leq j \leq r_1}$ are viewed as column vectors. Moreover, we require that $a_\alpha \neq 0$ for some open chart $\Omega \subset M$ and some choice of multi-index α with $\|\alpha\| = l$.

Let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be a \mathbb{C} -linear differential operator of order l from E_1 to E_2 . The *principal symbol*, or simply the *symbol*, of P is the map

$$\sigma_P : T^*M \rightarrow \text{Hom}(E_1, E_2) \quad (x, \xi) \mapsto \sum_{\|\alpha\|=l} a_\alpha(x) \xi^\alpha.$$

We say that P is *elliptic* if $\sigma_P(x, \xi) \in \text{Hom}((E_1)_x, (E_2)_x)$ is an isomorphism for all $x \in M$ and $0 \neq \xi \in T_x^*M$. Note that, if $u \in \Gamma(M, E_1)$ and $f \in C^\infty(M)$, with $f(x) = 0$ then

$$P(f^l u)(x) = l! \sigma_P(x, df(x))(u(x)).$$

Therefore, we observe that P is elliptic if and only if for all $x \in M$, $u \in \Gamma(M, E_1)$ and $f \in C^\infty(M)$ such that $u(x) \neq 0$, $f(x) = 0$ and $df(x) \neq 0$ we have

$$P(f^l u)(x) \neq 0.$$

Let (M, g) be an oriented Riemannian manifold of dimension m . The metric g induces the standard Riemannian volume form, given locally by

$$\text{Vol}(x) = |\det g_{ij}(x)|^{\frac{1}{2}} dx^1 \wedge \dots \wedge dx^m,$$

where $g(x) = \sum g_{ij}(x) dx^i \otimes dx^j$ for local coordinates x^1, \dots, x^m compatible with the orientation. Likewise, we also obtain the standard Riemannian measure from g .

Let E be a \mathbb{C} -vector bundle over M , and take a Hermitian metric h over E , *i.e.*, a smooth section of Hermitian inner products on the fibres. The couple (E, h) will be called a *Hermitian vector bundle*. The separable Banach space $L^p E, p \geq 1$, is then given by (equivalence classes of almost everywhere equal) possibly non-continuous global sections u of E , with measurable coefficients and finite L^p norm.

$$\|u\|_{L^p} := \left(\int_M |u(x)|^p \text{Vol}(x) \right)^{\frac{1}{p}} < +\infty,$$

where $|\cdot| = h(\cdot, \cdot)^{\frac{1}{2}}$. The space $L^p E$ can be seen as the completion, with respect to the L^p norm, of $\Gamma_0(M, E)$, the space of smooth sections with compact support. We

denote by $L^p_{loc}E$ the space of global sections u of E such that $u\chi_K \in L^pE$ for every compact $K \subseteq M$, where $\chi_K = 1$ on K and $\chi_K = 0$ otherwise. For $p = 2$, we denote the corresponding L^2 inner product by

$$\langle u, v \rangle := \int_M h(u(x), v(x)) \text{Vol}(x).$$

The space L^2E together with $\langle \cdot, \cdot \rangle$ is a separable Hilbert space. Denote by $\|\cdot\|$ the L^2 norm $\|\cdot\|_{L^2}$.

Let $(E_1, h_1), (E_2, h_2)$ be Hermitian vector bundles, and let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be a differential operator. The L^2 formal adjoint, or simply the formal adjoint,

$$P^* : \Gamma(M, E_2) \rightarrow \Gamma(M, E_1)$$

of P is defined such that it satisfies the property

$$\langle Pu, v \rangle = \langle u, P^*v \rangle$$

for any pair of smooth sections $u \in \Gamma(M, E_1)$ and $v \in \Gamma(M, E_2)$ with $\text{supp } u \cap \text{supp } v$ compactly contained in M . We remark that the L^2 formal adjoint P^* is a differential operator, it always exists and it is unique (see, e.g., [15, Chapter VI, Definition 1.5]). It follows that $P^{**} = P$ and, if Q is another differential operator, $(QP)^* = P^*Q^*$. An operator $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_1)$ is called *formally self-adjoint* if $P = P^*$.

As a generalisation of the notion of an elliptic differential operator, we recall the definition of an elliptic complex. We start by considering a complex of differential operators. Take E_j to be a sequence of \mathbb{C} -vector bundles over M and take D_j to be a sequence of \mathbb{C} -linear differential operators of order k_j

$$\dots \longrightarrow \Gamma(M, E_j) \xrightarrow{D_j} \Gamma(M, E_{j+1}) \xrightarrow{D_{j+1}} \Gamma(M, E_{j+2}) \longrightarrow \dots \tag{1}$$

such that $D_j \circ D_{j+1} = 0$ for all j . Given Hermitian metrics h_j on E_j , this complex of differential operators has naturally associated, formally self-adjoint operators

$$\Delta_j := D_j^*D_j + D_{j-1}D_{j-1}^* : \Gamma(M, E_j) \longrightarrow \Gamma(M, E_{j+1}) \tag{2}$$

of order $2 \max(k_j, k_{j-1})$ which are called Laplacians.

We can view the principal symbol of the operator D_j as a map

$$\sigma_{D_j} : \pi^*E_j \rightarrow \pi^*E_{j+1},$$

where $\pi : T^*M \setminus \{\text{zero section}\} \rightarrow M$ and π^*E_j denotes the pullback bundle of E_j . In this way, we say that (1) is an *elliptic complex* if the induced sequence of symbols

$$\cdots \longrightarrow \pi^* E_j \xrightarrow{\sigma_{D_j}} \pi^* E_{j+1} \xrightarrow{\sigma_{D_{j+1}}} \pi^* E_{j+2} \longrightarrow \cdots$$

is exact, *i.e.*, if $\text{im } \sigma_{D_j} = \ker \sigma_{D_{j+1}}$ for all j . Note that in general, even for elliptic complexes, Δ_j is not an elliptic operator.

There is, however, a different formally self-adjoint operator naturally associated to the elliptic complex, which is always elliptic (see [43, Section I]). Focusing just on two consecutive maps in the complex

$$\Gamma(M, E_1) \xrightarrow{D_1} \Gamma(M, E_2) \xrightarrow{D_2} \Gamma(M, E_3),$$

For $j = 1, 2$ we set

$$l_j := \frac{\text{lcm}(k_1, k_2)}{k_j}, \quad r := k_1 l_1 = k_2 l_2 = \text{lcm}(k_1, k_2),$$

and define the operator

$$\square := (D_1 D_1^*)^{l_1} + (D_2^* D_2)^{l_2} \tag{3}$$

of order $2r$, which is clearly formally self-adjoint. To see that \square is elliptic, one can argue similarly to [31, Lemma 9.4.2].

Finally, we state the following result about elliptic regularity, whose proof follows, *e.g.*, from [27, Lemma 1.1.17]. Let $(E_1, h_1), (E_2, h_2)$ be Hermitian vector bundles, and let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be a differential operator. We say that the section u is a *weak solution* of $Pu = v$ if $u \in L^1_{loc} E_1, v \in L^1_{loc} E_2$ and

$$\langle u, P^* w \rangle = \langle v, w \rangle \quad \forall w \in \Gamma_0(M, E_2).$$

Theorem 2.1 (*Elliptic regularity*). *Given an oriented Riemannian manifold (M, g) , let $(E_1, h_1), (E_2, h_2)$ be a pair of Hermitian vector bundles over M and let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be an elliptic differential operator. If $u \in L^1_{loc} E_1, u$ is a weak solution of $Pu = v$ and v is smooth, then u must be smooth.*

3. Elliptic complexes on complex manifolds

Good references for the content of this section are [15] and [39]. Let M be a complex manifold of complex dimension n . We will use $\Lambda^r M, \Lambda^r_{\mathbb{C}} M$ and $\Lambda^{p,q} M$ to denote the bundles of real-valued r -forms, complex-valued r -forms, and (p, q) -forms, respectively. When the manifold is obvious from context, we will omit it and simply write $\Lambda^r, \Lambda^r_{\mathbb{C}}$ and $\Lambda^{p,q}$. The spaces of smooth sections of these bundles are then given by

$$A^r(M) := \Gamma(M, \Lambda^r M), \quad A^r_{\mathbb{C}}(M) := \Gamma(M, \Lambda^r_{\mathbb{C}} M), \quad A^{p,q}(M) := \Gamma(M, \Lambda^{p,q} M),$$

or simply $A^r, A^r_{\mathbb{C}}$ and $A^{p,q}$.

The exterior derivative on $A^{p,q}$ splits into two bidegree components

$$d = \partial + \bar{\partial},$$

where

$$\partial : A^{p,q} \rightarrow A^{p+1,q}, \quad \bar{\partial} : A^{p,q} \rightarrow A^{p,q+1},$$

therefore the relation $d^2 = 0$ immediately implies

$$\partial^2 = 0, \quad \bar{\partial}^2 = 0, \quad \partial\bar{\partial} + \bar{\partial}\partial = 0.$$

Thanks to these relations there are complexes of differential forms which are naturally associated to a complex manifold.

The first complex we are interested in is the Dolbeault complex, given by

$$\dots \rightarrow A^{p,q-1} \xrightarrow{\bar{\partial}} A^{p,q} \xrightarrow{\bar{\partial}} A^{p,q+1} \rightarrow \dots$$

for all $0 \leq p, q \leq n$, along with the associated Dolbeault cohomology

$$H_{\bar{\partial}}^{p,q} := \frac{\ker \bar{\partial} \cap A^{p,q}}{\text{im } \bar{\partial} \cap A^{p,q}}.$$

There is an analogous complex for the map ∂ , with the associated cohomology $H_{\partial}^{p,q}$.

The second complex we are interested in is the Aeppli-Bott-Chern complex, given by

$$\begin{array}{c} \dots \\ \downarrow \\ A^{p-1,q-2} \oplus A^{p-2,q-1} \\ \downarrow \bar{\partial} \oplus \partial \\ A^{p-1,q-1} \\ \downarrow \partial \bar{\partial} \\ A^{p,q} \\ \downarrow \partial + \bar{\partial} \\ A^{p+1,q} \oplus A^{p,q+1} \\ \downarrow \\ \dots \end{array} \tag{4}$$

for all $0 \leq p, q \leq n$. The complete definition of the ABC complex is given in Section 10 (cf. [15, Chapter VI, Section 12.1]). The Bott-Chern and Aeppli cohomologies

$$H_{BC}^{p,q} := \frac{\ker \partial \cap \ker \bar{\partial} \cap A^{p,q}}{\text{im } \partial \bar{\partial} \cap A^{p,q}}, \quad H_A^{p-1,q-1} := \frac{\ker \partial \bar{\partial} \cap A^{p-1,q-1}}{\text{im } \bar{\partial} \oplus \partial \cap A^{p-1,q-1}}$$

can be obtained from this complex.

There are well-defined maps between these cohomology spaces, induced by the identity on the representatives of the cohomology classes. Specifically we have the following commutative diagram of maps

$$\begin{array}{ccccc}
 & & H_{BC}^{\bullet, \bullet} & & \\
 & \swarrow & \downarrow & \searrow & \\
 H_{\partial}^{\bullet, \bullet} & & H_{dR}^{\bullet, \bullet} & & H_{\bar{\partial}}^{\bullet, \bullet} \\
 & \searrow & \downarrow & \swarrow & \\
 & & H_A^{\bullet, \bullet} & &
 \end{array} \tag{5}$$

where H_{dR}^k is the de Rham cohomology. Moreover, by [13, Lemma 5.15, Remark 5.16], these maps are all isomorphisms if and only if one of the following equivalent conditions holds on $A_{\mathbb{C}}^{\bullet}$:

- a) $\text{im } \partial\bar{\partial} = \ker \partial \cap \ker \bar{\partial} \cap \text{im } d$;
- b) $\text{im } \partial\bar{\partial} = \ker \partial \cap \text{im } \bar{\partial}$;
- c) $\text{im } \partial\bar{\partial} = \ker \partial \cap \ker \bar{\partial} \cap (\text{im } \partial + \text{im } \bar{\partial})$;
- d) $\ker \partial\bar{\partial} = \text{im } \partial + \text{im } \bar{\partial} + \ker d$;
- e) $\ker \partial\bar{\partial} = \text{im } \bar{\partial} + \ker \partial$;
- f) $\ker \partial\bar{\partial} = \text{im } \partial + \text{im } \bar{\partial} + (\ker \partial \cap \ker \bar{\partial})$.

Note that when condition a) is satisfied it is often said that the $\partial\bar{\partial}$ -Lemma holds.

Using linear algebra one can show that the Dolbeault and the Aeppli-Bott-Chern complexes are both elliptic, see [43, Section II] or [40, Lemma 2]. Therefore, once we fix a Hermitian metric for differential forms, there are naturally associated elliptic and formally self-adjoint operators as defined in (3).

Let (M, g) be a Hermitian manifold of complex dimension n , that is a complex manifold M endowed with a Hermitian metric g . Recall that a Hermitian metric on a complex manifold is a Riemannian metric g for which the complex structure J is an isometry, *i.e.*, $g(J\cdot, J\cdot) = g(\cdot, \cdot)$. We will generally denote by ω the fundamental $(1, 1)$ -form associated to the metric g , which is defined by $\omega(\cdot, \cdot) := g(J\cdot, \cdot)$. We will also typically denote by h the Hermitian extension of g on the complexified tangent bundle $T^{\mathbb{C}}M = TM \otimes \mathbb{C} = T^{1,0}M \oplus T^{0,1}M$, and by the same symbol g the \mathbb{C} -bilinear symmetric extension of g on $T^{\mathbb{C}}M$. Consequently, we have $h(u, v) = g(u, \bar{v})$ for all $u, v \in \Gamma(M, T^{1,0}M)$. Note that the standard Riemannian volume form can be computed by $\text{Vol} = \frac{\omega^n}{n!}$.

The Hermitian metric g extends to a Hermitian metric h on $\Lambda^{p,q}M$, defined pointwise as follows. At any point $x \in M$, we choose a basis $\{v_1, \dots, v_n\}$ of $T_x^{1,0}M$, such that

$h_x(v_i, v_j) = \delta_{ij}$. Take the dual basis $\{\alpha_1, \dots, \alpha_n\}$ of $\Lambda_x^{1,0}$ and define h_x on $\Lambda_x^{1,0}$ in such a way that $h_x(\alpha_i, \alpha_j) = \delta_{ij}$. We then extend the Hermitian metric to $\Lambda_x^{p,q}$ by setting

$$h_x(\alpha^{i_1} \wedge \dots \wedge \bar{\alpha}^{j_q}, \alpha^{k_1} \wedge \dots \wedge \bar{\alpha}^{h_q}) = \delta_{i_1 k_1} \dots \delta_{j_q h_q}.$$

Then, $L^2\Lambda^{p,q}$ is the space of (equivalence classes of almost everywhere equal) measurable (p, q) -forms φ , such that

$$\|\varphi\| := \left(\int_M h(\varphi, \varphi) \text{Vol} \right)^{\frac{1}{2}} < \infty.$$

Pairing $L^2\Lambda^{p,q}$ with the Hermitian inner product

$$\langle \varphi, \psi \rangle := \int_M h(\varphi, \psi) \text{Vol},$$

we get a Hilbert space. The space $L^2\Lambda^{p,q}$ can be also seen as the completion of $A_0^{p,q} := \Gamma_0(M, \Lambda^{p,q})$, the space of smooth (p, q) -forms with compact support, with respect to the norm $\|\cdot\|$.

The complex \mathbb{C} -linear Hodge operator $*$: $A^{p,q} \rightarrow A^{n-q, n-p}$ associated with the metric is defined by the equation

$$\alpha \wedge *\bar{\beta} = h(\alpha, \beta) \text{Vol}$$

for all $\alpha, \beta \in A^{p,q}$. We set also

$$\bar{\partial}^* := -*\partial* \quad \partial^* := -*\bar{\partial}*, \quad d^* := -*d*,$$

which are the L^2 -formal adjoints respectively of $\partial, \bar{\partial}, d$ by a direct application of the Stokes Theorem.

Note that the Laplacians defined in (2) and (3) coincide, and so there is a single elliptic and formally self-adjoint Laplacian associated to the Dolbeault complex,

$$\Delta_{\bar{\partial}} := \bar{\partial}\bar{\partial}^* + \bar{\partial}^*\bar{\partial},$$

known as the *Dolbeault Laplacian*. Similarly we define

$$\Delta_{\partial} := \partial\partial^* + \partial^*\partial, \quad \Delta_d := dd^* + d^*d,$$

which are respectively called ∂ -Laplacian and the *Hodge Laplacian*.

For the Aepli-Bott-Chern complex we obtain different Laplacians from (2) and (3). Note that the formal adjoint of the map $\partial + \bar{\partial}$ in (4) is $\partial^* \oplus \bar{\partial}^*$, while the formal adjoint of $\bar{\partial} \oplus \partial$ in (4) is $\bar{\partial}^* + \partial^*$. The operators

$$\begin{aligned} \Delta_{BC} &:= \partial\bar{\partial}\bar{\partial}^*\partial^* + (\partial^* \oplus \bar{\partial}^*)(\partial + \bar{\partial}) = \partial\bar{\partial}\bar{\partial}^*\partial^* + \partial^*\partial + \bar{\partial}^*\bar{\partial}, \\ \Delta_A &:= \bar{\partial}^*\partial^*\partial\bar{\partial} + (\bar{\partial} \oplus \partial)(\partial^* + \bar{\partial}^*) = \bar{\partial}^*\partial^*\partial\bar{\partial} + \partial\partial^* + \bar{\partial}\bar{\partial}^*, \end{aligned}$$

acting respectively on $A^{p,q}$ and $A^{p-1,q-1}$, are the Laplacians given by (2). Recall that they are formally self-adjoint but, as noted in [39], they are not elliptic. Whereas, the operators

$$\begin{aligned} \square_{BC} &:= \partial\bar{\partial}\bar{\partial}^*\partial^* + ((\partial^* \oplus \bar{\partial}^*)(\partial + \bar{\partial}))^2 = \partial\bar{\partial}\bar{\partial}^*\partial^* + (\partial^*\partial + \bar{\partial}^*\bar{\partial})^2, \\ \square_A &:= \bar{\partial}^*\partial^*\partial\bar{\partial} + ((\bar{\partial} \oplus \partial)(\partial^* + \bar{\partial}^*))^2 = \bar{\partial}^*\partial^*\partial\bar{\partial} + (\partial\partial^* + \bar{\partial}\bar{\partial}^*)^2, \end{aligned}$$

acting respectively on $A^{p,q}$ and $A^{p-1,q-1}$, are the Laplacians given by (3) and are both formally self-adjoint and elliptic [43, Section 2].

Moreover, Kodaira and Spencer [26] have introduced the following fourth order elliptic and formally self-adjoint differential operators usually referred to as the *Bott-Chern* and *Aeppli Laplacians*, which are defined respectively as

$$\begin{aligned} \tilde{\Delta}_{BC} &:= \partial\bar{\partial}\bar{\partial}^*\partial^* + \bar{\partial}^*\partial^*\partial\bar{\partial} + \partial^*\bar{\partial}\bar{\partial}^*\partial + \bar{\partial}^*\partial\partial^*\bar{\partial} + \partial^*\partial + \bar{\partial}^*\bar{\partial}, \\ \tilde{\Delta}_A &:= \partial\bar{\partial}\bar{\partial}^*\partial^* + \bar{\partial}^*\partial^*\partial\bar{\partial} + \partial\bar{\partial}^*\bar{\partial}\partial^* + \bar{\partial}\bar{\partial}^*\partial\bar{\partial}^* + \partial\partial^* + \bar{\partial}\bar{\partial}^*. \end{aligned}$$

All these Laplacians are linked by the following duality relations

$$\begin{aligned} *\Delta_A &= \Delta_{BC}*, & *\Delta_{BC} &= \Delta_A*, \\ *\tilde{\Delta}_A &= \tilde{\Delta}_{BC}*, & *\tilde{\Delta}_{BC} &= \tilde{\Delta}_A*, \\ *\square_A &= \square_{BC}*, & *\square_{BC} &= \square_A*. \end{aligned} \tag{6}$$

Let (M, g) be a compact Hermitian manifold. It is then straightforward to see

$$\begin{aligned} \ker \Delta_{\bar{\partial}} &= \ker \bar{\partial} \cap \ker \bar{\partial}^*, & \ker \Delta_{\partial} &= \ker \partial \cap \ker \partial^*, & \ker \Delta_d &= \ker d \cap \ker d^*, \\ \ker \Delta_{BC} &= \ker \tilde{\Delta}_{BC} = \ker \square_{BC} = \ker \bar{\partial}^*\partial^* \cap \ker \partial \cap \ker \bar{\partial}, \\ \ker \Delta_A &= \ker \tilde{\Delta}_A = \ker \square_A = \ker \partial\bar{\partial} \cap \ker \partial^* \cap \ker \bar{\partial}^*. \end{aligned}$$

Notice that the kernels of the Bott-Chern Laplacians $\Delta_{BC}, \tilde{\Delta}_{BC}$ and \square_{BC} all coincide, and likewise for the Aeppli Laplacian. The respective spaces of harmonic forms will be denoted by

$$\mathcal{H}_d^k := \ker \Delta_d \cap A^k, \quad \mathcal{H}_\delta^{p,q} := \ker \Delta_\delta \cap A^{p,q},$$

for $\delta \in \{\partial, \bar{\partial}, BC, A\}$, as a consequence of Hodge theory they are finite dimensional and isomorphic to the respective cohomology spaces

$$\mathcal{H}_d^k \simeq H_{dR}^k, \quad \mathcal{H}_\delta^{p,q} \simeq H_\delta^{p,q}.$$

Their dimensions will be indicated by

$$b^k := \dim_{\mathbb{R}} H_{dR}^k, \quad h_{\delta}^{p,q} := \dim_{\mathbb{C}} H_{\delta}^{p,q}.$$

If (M, g) is a Kähler manifold (i.e., $d\omega = 0$, where ω is the fundamental form), it is well known that the second order Laplacians coincide up to a factor,

$$\Delta_d = 2\Delta_{\partial} = 2\Delta_{\bar{\partial}}. \tag{7}$$

Furthermore, using the Kähler identities, see, e.g., [15, Chapter VI, Theorem 6.4], we know that ∂ and $\bar{\partial}^*$ anticommute, as do ∂^* and $\bar{\partial}$. This allows us to write $\tilde{\Delta}_{BC}$ and $\tilde{\Delta}_A$ in a more concise form,

$$\tilde{\Delta}_{BC} = \Delta_{\bar{\partial}}\Delta_{\bar{\partial}} + \partial^*\partial + \bar{\partial}^*\bar{\partial}, \tag{8}$$

$$\tilde{\Delta}_A = \Delta_{\bar{\partial}}\Delta_{\bar{\partial}} + \partial\partial^* + \bar{\partial}\bar{\partial}^*. \tag{9}$$

As a consequence, if M is compact, all the spaces of harmonic forms coincide

$$\mathcal{H}_d^k \otimes \mathbb{C} \cap A^{p,q} = \mathcal{H}_{\partial}^{p,q} = \mathcal{H}_{\bar{\partial}}^{p,q} = \mathcal{H}_{BC}^{p,q} = \mathcal{H}_A^{p,q}.$$

In particular, this implies that the $\partial\bar{\partial}$ -Lemma holds, see [13, Lemma 5.11, Remark 5.14].

4. Unbounded operators on Hilbert spaces

We recall some concepts from the theory of unbounded operators on Hilbert spaces. For a complete treatment we refer to [18, Chapter 1] or [36, Chapter 8]. If $\mathcal{H}_1, \mathcal{H}_2$ are Hilbert spaces, a linear operator $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ is a linear function $P : \mathcal{D}(P) \rightarrow \mathcal{H}_2$, defined on a domain $\mathcal{D}(P) \subseteq \mathcal{H}_1$. The graph of a linear operator $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ is the subspace of $\mathcal{H}_1 \times \mathcal{H}_2$ given by $\text{Gr}(P) := \{(x, Px) \in \mathcal{H}_1 \times \mathcal{H}_2 \mid x \in \mathcal{D}(P)\}$. A linear operator is closed if its graph is a closed subspace. By the closed graph theorem, a closed linear operator which is defined everywhere on \mathcal{H}_1 is automatically bounded, therefore when dealing with an unbounded operator we must always keep track of its domain.

The kernel of a linear operator $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ with domain $\mathcal{D}(P)$ is the subspace $\ker P := \{x \in \mathcal{D}(P) \mid Px = 0\}$, while its image is $\text{im } P := P(\mathcal{D}(P))$. If P is closed, then its kernel is a closed subspace.

An extension of P is a linear operator P' such that $\mathcal{D}(P) \subseteq \mathcal{D}(P')$ and $Px = P'x$ for every $x \in \mathcal{D}(P)$, i.e., $\text{Gr}(P) \subseteq \text{Gr}(P')$; in this case we will write $P \subseteq P'$. A linear operator P is closable if it admits a closed extension. Every closable operator P has a smallest closed extension, denoted by \bar{P} , which is called its closure, and is given by the intersection of the graphs of all closed extensions. Note that a linear subspace G of $\mathcal{H}_1 \times \mathcal{H}_2$ is a graph of a linear operator so long as $(0, y) \notin G$ for all non-zero $y \in \mathcal{H}_2$. If P is closable then $\text{Gr}(\bar{P}) = \overline{\text{Gr}(P)}$, [36, p. 250].

If $\mathcal{D}(P)$ is dense in \mathcal{H}_1 , then we say that the linear operator P is *densely defined*, and we can define the *Hilbert adjoint*, or simply the *adjoint*, of P , indicated by $P^t : \mathcal{H}_2 \rightarrow \mathcal{H}_1$. Its domain is

$$\mathcal{D}(P^t) := \{y \in \mathcal{H}_2 \mid x \mapsto \langle Px, y \rangle_2 \text{ is bounded on } \mathcal{D}(P)\},$$

where $\langle \cdot, \cdot \rangle_i$ denotes the Hermitian inner product of the Hilbert space \mathcal{H}_i . The adjoint is then defined by the relation

$$\langle Px, y \rangle_2 = \langle x, P^t y \rangle_1 \quad \forall x \in \mathcal{D}(P) \quad \forall y \in \mathcal{D}(P^t).$$

Indeed, if $y \in \mathcal{D}(P^t)$, then the map $x \mapsto \langle Px, y \rangle_2$ can be uniquely extended from a densely defined function on $\mathcal{D}(P)$ to a bounded linear function on all of \mathcal{H}_1 . By the Riesz representation theorem, there exists a unique $z \in \mathcal{H}_1$ such that $\langle Px, y \rangle_2 = \langle x, z \rangle_1$, for all $x \in \mathcal{D}(P)$. We then define $P^t y := z$.

We can verify that $\text{Gr}(P^t) = (\text{Gr}(-P))^\perp$ in $\mathcal{H}_1 \times \mathcal{H}_2$, thus this definition makes P^t a closed linear operator. If P is closed, this implies that every pair $(u, v) \in \mathcal{H}_1 \times \mathcal{H}_2$ can be written as the sum of elements in $\text{Gr}(P^t)$ and $\text{Gr}(-P)$,

$$(u, v) = (x, -Px) + (P^t y, y), \quad x \in \mathcal{D}(P), \quad y \in \mathcal{D}(P^t).$$

If $u = 0$, then we have

$$x + P^t y = 0, \quad v = y - Px = y + PP^t y, \quad \langle v, y \rangle_2 = \|y\|_2^2 + \|P^t y\|_1^2,$$

where $\|\cdot\|_i := \langle \cdot, \cdot \rangle_i^{\frac{1}{2}}$. If $v \in \mathcal{D}(P^t)^\perp$, we get $\langle v, y \rangle_2 = 0$, thus $y = 0$ and $v = 0$. Therefore, P^t is densely defined and so we can define P^{tt} . In fact, we can verify that $P^{tt} = P$.

The above discussion (see also [36, Theorem VIII.1]) implies the following result.

Lemma 4.1. *Let $\mathcal{H}_1, \mathcal{H}_2$ be two complex Hilbert spaces and $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ be a linear densely defined operator. If P is closed, then its adjoint P^t is closed and densely defined, $P^{tt} = P$ and*

$$\ker P^t = \text{im } P^\perp \quad \ker P^\perp = \overline{\text{im } P^t}.$$

Moreover P is closable iff $\mathcal{D}(P^t)$ is dense, in which case $(\overline{P})^t = P^t$ and $\overline{P} = P^{tt}$.

Given two densely defined linear operators P, Q , it holds that

$$\begin{aligned} P \subseteq Q &\iff Q^t \subseteq P^t, \\ \text{im } P \subseteq \ker Q &\iff \text{im } Q^t \subseteq \ker P^t. \end{aligned}$$

It follows that any Hilbert complex

$$\mathcal{H}_1 \xrightarrow{P} \mathcal{H}_2 \xrightarrow{Q} \mathcal{H}_3,$$

i.e., any sequence of closed and densely defined linear operators P, Q , such that $\text{im } P \subseteq \ker Q$, has an associated *dual Hilbert complex*

$$\mathcal{H}_1 \xleftarrow{P^t} \mathcal{H}_2 \xleftarrow{Q^t} \mathcal{H}_3.$$

Another fundamental property of Hilbert complexes is the following orthogonal decomposition.

Theorem 4.2. *Let $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ and $Q : \mathcal{H}_2 \rightarrow \mathcal{H}_3$ be closed and densely defined linear operators between Hilbert spaces. If $\text{im } P \subseteq \ker Q$, then we have the orthogonal decompositions*

$$\begin{aligned} \mathcal{H}_2 &= \ker Q \cap \ker P^t \oplus \overline{\text{im } P} \oplus \overline{\text{im } Q^t}, \\ \ker Q &= \ker Q \cap \ker P^t \oplus \overline{\text{im } P}, \\ \ker P^t &= \ker Q \cap \ker P^t \oplus \overline{\text{im } Q^t}. \end{aligned}$$

Proof. By Lemma 4.1, the Hilbert space \mathcal{H}_2 has an orthogonal decomposition

$$\mathcal{H}_2 = \ker Q \oplus \overline{\text{im } Q^t},$$

and analogously

$$\mathcal{H}_2 = \ker P^t \oplus \overline{\text{im } P}.$$

Furthermore, since $\ker Q$ is a closed subspace of \mathcal{H}_2 , it is itself a Hilbert space with $\overline{\text{im } P} \subseteq \ker Q$. Again by Lemma 4.1, we see that $\ker Q$ has an orthogonal decomposition

$$\ker Q = \ker Q \cap \ker P^t \oplus \overline{\text{im } P}.$$

The last decomposition follows similarly. \square

The following basic lemma characterises the closure of the image of a closed and densely defined operator. We include a proof for the convenience of the reader.

Lemma 4.3. *Let $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ be a closed and densely defined linear operator between Hilbert spaces. The following conditions are equivalent:*

- a) $\text{im } P$ is closed;
- b) $\exists C > 0$ s.t. $\|x\|_1 \leq C\|Px\|_2$ for all $x \in \mathcal{D}(P) \cap \overline{\text{im } P^t}$;
- c) $\text{im } P^t$ is closed;

d) $\exists C > 0$ s.t. $\|y\|_2 \leq C\|P^t y\|_1$ for all $y \in \mathcal{D}(P^t) \cap \overline{\text{im } P}$.

Proof. By Lemma 4.1, the Hilbert space \mathcal{H}_1 can be orthogonally decomposed as

$$\mathcal{H}_1 = \ker P \oplus \overline{\text{im } P^t}, \tag{10}$$

which implies that the operator $P|_{\mathcal{D}(P) \cap \overline{\text{im } P^t}} : \mathcal{D}(P) \cap \overline{\text{im } P^t} \rightarrow \text{im } P$ is bijective, therefore its inverse is well-defined and one can check that it is also a closed operator. If a) holds, then by the closed graph theorem $(P|_{\mathcal{D}(P) \cap \overline{\text{im } P^t}})^{-1}$ is continuous, proving b). Conversely, if b) holds, then using (10) we can show that any Cauchy sequence in $\text{im } P$ converges in $\text{im } P$, and thus we obtain a). Similarly c) is equivalent to d). To prove that b) implies d), note that by b)

$$|\langle y, Px \rangle_2| = |\langle P^t y, x \rangle_1| \leq C\|P^t y\|_1 \|Px\|_2$$

for all $y \in \mathcal{D}(P^t)$ and $x \in \mathcal{D}(P) \cap \overline{\text{im } P^t}$. Therefore

$$|\langle y, z \rangle_2| \leq C\|P^t y\|_1 \|z\|_2$$

for all $y \in \mathcal{D}(P^t)$ and $z \in \overline{\text{im } P}$, again using (10). Choosing $z = y$ then implies d). Since $P^{tt} = P$, a similar argument shows that d) implies b). \square

Let \mathcal{H} be a Hilbert space. A densely defined linear operator $P : \mathcal{H} \rightarrow \mathcal{H}$ is called *symmetric* if $\langle Px, y \rangle = \langle x, Py \rangle$ whenever $x, y \in \mathcal{D}(P)$, or equivalently if $P \subseteq P^t$. It is called *self-adjoint* if it is symmetric and $\mathcal{D}(P) = \mathcal{D}(P^t)$, or equivalently if $P = P^t$ (with equality of domains). A symmetric operator P is always closable since its adjoint is a closed extension, therefore $\overline{P} = P^{tt}$ and so $P \subseteq P^{tt} \subseteq P^t$. A linear operator P is called *essentially self-adjoint* if it has a unique self-adjoint extension. Equivalently, P is essentially self-adjoint if \overline{P} is a self-adjoint operator, [36, p. 256]. An operator is called *positive* if $\langle Px, x \rangle \geq 0$ whenever $x \in \mathcal{D}(P)$.

The following Theorems describe a method for building self-adjoint operators.

Theorem 4.4 ([37, Theorem X.25] or [18, Theorem 2.3]). *Let P be a closed and densely defined linear operator on a Hilbert space \mathcal{H} . Then the operator $P^t P$ defined by $(P^t P)x = P^t(Px)$ on the domain*

$$\mathcal{D}(P^t P) := \{x \in \mathcal{D}(P) \mid Px \in \mathcal{D}(P^t)\}$$

is positive and self-adjoint.

Theorem 4.5 ([18, Theorem 4.1]). *Let P, Q be positive and self-adjoint operators on a Hilbert space \mathcal{H} . Assume that $\mathcal{D}(P) \cap \mathcal{D}(Q)$ is dense in \mathcal{H} . Then the operator $P + Q$ defined by $(P + Q)x = Px + Qx$ on the domain*

$$\mathcal{D}(P + Q) := \mathcal{D}(P) \cap \mathcal{D}(Q)$$

is positive and self-adjoint.

By Theorems 4.4 and 4.5, any Hilbert complex

$$\mathcal{H}_1 \xrightarrow{P} \mathcal{H}_2 \xrightarrow{Q} \mathcal{H}_3$$

has a naturally associated Laplacian, given by

$$\Delta := PP^t + Q^tQ, \tag{11}$$

which is positive and self-adjoint. Note that the Laplacian of a Hilbert complex coincides with the Laplacian of its dual Hilbert complex. We will need later on the following observation: being Δ self adjoint, we can define Δ^2 by Theorem 4.4, and its domain turns out to satisfy

$$\mathcal{D}(\Delta^2) = \mathcal{D}((PP^t)^2) \cap \mathcal{D}((Q^tQ)^2). \tag{12}$$

Finally, it is easy to check that the kernel of the Laplacian is characterised by

$$\ker \Delta = \ker Q \cap \ker P^t.$$

More generally, we have the following result.

Lemma 4.6. *Let $P_j : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ be closed and densely defined linear operators for $j = 1, \dots, n$, $n \in \mathbb{N}$, and define $\square := \sum_{j=1}^n P_j^t P_j$ with the domain given by Theorems 4.4 and 4.5 (we ask that $\bigcap_{j=1}^n \mathcal{D}(P_j^t P_j)$ is dense in \mathcal{H}_1). Then*

$$\ker \square = \bigcap_{j=1}^n \ker P_j.$$

4.1. Spectral gap of the Laplacian of a Hilbert complex

We end this section by analysing some characterisations of the spectral gap condition of any positive self-adjoint operator and in particular of the Laplacian of a given Hilbert complex.

We say that $\lambda \in \mathbb{C}$ belongs to the *spectrum* $\sigma(P)$ of an unbounded linear operator $P : \mathcal{H} \rightarrow \mathcal{H}$ if there exists no bounded linear operator $B : \mathcal{H} \rightarrow \mathcal{H}$ such that

- 1) $B(P - \lambda I)x = x$ for all $x \in \mathcal{D}(P)$,
- 2) $Bx \in \mathcal{D}(P)$ and $(P - \lambda I)Bx = x$ for all $x \in \mathcal{H}$;

in other words, if $P - \lambda I$ has no bounded inverse. By [37, Theorem X.1], the spectrum of a self-adjoint operator is a subset of the real axis, and the spectrum of a positive self-adjoint operator is a subset of the non-negative real axis. Furthermore, we say that a positive self-adjoint operator P has a *spectral gap* if it has the property $\inf(\sigma(P) \setminus \{0\}) = C > 0$, or equivalently if $\sigma(P) \subseteq \{0\} \cup [C, +\infty)$ with $C > 0$.

We now state the Spectral Theorem for unbounded self-adjoint operators. We will introduce very briefly the notion of a direct integral.

Let (X, μ) be a σ -finite measure space and $\{\mathbf{H}_\lambda\}_{\lambda \in X}$ be a collection of separable Hilbert spaces and a *measurability structure*, see [22, Definition 7.18] for more details. Denote by $\langle \cdot, \cdot \rangle_\lambda$ and $\|\cdot\|_\lambda$ the inner product and the norm on \mathbf{H}_λ . A *section* s is a function $X \rightarrow \bigcup_{\lambda \in X} \mathbf{H}_\lambda$ satisfying $s(\lambda) \in \mathbf{H}_\lambda$ for all $\lambda \in X$. The *direct integral*

$$\int_X^\oplus \mathbf{H}_\lambda d\mu(\lambda)$$

is the Hilbert space of classes of a.e. equal *measurable* sections s with finite norm $\|s\| < +\infty$, where $\|s\| := \langle s, s \rangle^{\frac{1}{2}}$ and

$$\langle s_1, s_2 \rangle := \int_X \langle s_1(\lambda), s_2(\lambda) \rangle_\lambda d\mu(\lambda).$$

Theorem 4.7 (Spectral Theorem). *Let P be a self-adjoint operator on a separable Hilbert space \mathcal{H} . Then there is a σ -finite measure μ on the spectrum $\sigma(P)$, along with a unitary map (a linear bijection preserving the inner products)*

$$U : \mathcal{H} \rightarrow \int_{\sigma(P)}^\oplus \mathbf{H}_\lambda d\mu(\lambda)$$

such that

$$U(\mathcal{D}(P)) := \left\{ s \in \int_{\sigma(P)}^\oplus \mathbf{H}_\lambda d\mu(\lambda) : \int_{\sigma(P)} \|\lambda s(\lambda)\|_\lambda^2 d\mu(\lambda) < +\infty \right\}$$

and, for all $s \in U(\mathcal{D}(P))$ and $\lambda \in \sigma(P)$,

$$(UPU^{-1}s)(\lambda) = \lambda s(\lambda). \tag{13}$$

In this case, we say that P is isomorphic to the multiplication (operator) by λ , *i.e.*, the identity function on $\sigma(P)$. We refer to [22, Theorem 10.09, Section 10.4] for the proof. The point of the Spectral Theorem is that any question about a single self-adjoint

operator is a question about a function. For example, we can apply measurable functions to self-adjoint operators as follows. If φ is a complex Borel measurable function on the reals, then $\varphi(P)$ is defined as the operator which is isomorphic (by the same unitary operator) to multiplication by $\varphi(\lambda)$. E.g., the integer powers P^k are well-defined for $k \in \mathbb{Z}$ and this definition of P^2 coincides with the one given by applying Theorem 4.4.

With the help of the Spectral Theorem, we can prove the following characterisation of the spectral gap property.

Lemma 4.8. *Let P be a positive self-adjoint operator on a separable Hilbert space \mathcal{H} . Then the following conditions are equivalent:*

- a) P has a spectral gap;
- b) $\exists C > 0 \langle x, Px \rangle \geq C \langle x, x \rangle$ for all $x \in \mathcal{D}(P) \cap (\ker P)^\perp$;
- c) $\text{im } P$ is closed.

Proof. By Lemma 4.1 and Lemma 4.3 it is immediate to see b) is equivalent to c). We now prove that a) is equivalent to b). For all $x \in \mathcal{D}(P)$, by the Spectral Theorem let $Ux = s$, so that

$$\langle x, Px \rangle = \int_{\sigma(P)} \langle s(\lambda), \lambda s(\lambda) \rangle_\lambda d\mu(\lambda) = \int_{\sigma(P) \setminus \{0\}} \lambda \|s(\lambda)\|_\lambda^2 d\mu(\lambda).$$

Notice that $y \in \ker P$ iff, given $r = Uy$, $\lambda r(\lambda) = 0$ a.e., iff $r(\lambda) = 0$ a.e. for $\lambda \neq 0$. If $x \in \mathcal{D}(P) \cap (\ker P)^\perp$, then $s(0) = 0$, so that

$$\langle x, x \rangle = \int_{\sigma(P)} \langle s(\lambda), s(\lambda) \rangle_\lambda d\mu(\lambda) = \int_{\sigma(P) \setminus \{0\}} \|s(\lambda)\|_\lambda^2 d\mu(\lambda)$$

Therefore $\langle x, Px \rangle \geq C \langle x, x \rangle$ for all $x \in \mathcal{D}(P) \cap (\ker P)^\perp$ iff $\lambda \geq C$ a.e. in $\sigma(P) \setminus \{0\}$, iff $\sigma(P) \subseteq \{0\} \cup [C, +\infty)$. \square

If the positive self-adjoint operator is the Laplacian associated to a Hilbert complex, then we obtain other characterisations of the spectral gap condition.

Lemma 4.9. *Let $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ and $Q : \mathcal{H}_2 \rightarrow \mathcal{H}_3$ be closed and densely defined linear operators between separable Hilbert spaces satisfying $\text{im } P \subseteq \ker Q$. Define the positive self-adjoint operator*

$$\Delta = Q^t Q + P P^t,$$

with the domain given by Theorems 4.4 and 4.5 (we ask that $\mathcal{D}(Q^t Q) \cap \mathcal{D}(P P^t)$ is dense in \mathcal{H}_2). Then the following conditions are equivalent:

- a) Δ has a spectral gap;
- b) $\exists C > 0$ $C\|x\|_2^2 \leq \|P^t x\|_1^2 + \|Qx\|_3^2$ for all $x \in \mathcal{D}(P^t) \cap \mathcal{D}(Q) \cap (\ker \Delta)^\perp$;
- c) $\text{im } Q^t$ and $\text{im } P$ are closed.

Proof. Since $\ker \Delta = \ker P^t \cap \ker Q$, then $(\ker \Delta)^\perp = \overline{\text{im } P} \oplus \overline{\text{im } Q^t}$ by Lemma 4.1. Therefore c) implies b) applying Lemma 4.3. Moreover, b) implies a) thanks to the second characterisation in Lemma 4.8. To prove that a) implies c), by Lemma 4.8 we know that $\text{im } \Delta$ is closed, therefore by Lemma 4.1 we can decompose

$$\mathcal{H}_2 = \ker \Delta \oplus \text{im } \Delta \subseteq \ker \Delta \oplus \text{im } P \oplus \text{im } Q^t \subseteq \mathcal{H}_2,$$

thus

$$\mathcal{H}_2 = \ker \Delta \oplus \text{im } P \oplus \text{im } Q^t.$$

The orthogonality of the previous decompositions follows from Lemma 4.1 and by the assumption $\text{im } P \subseteq \ker Q$. The closure of $\text{im } P$ and $\text{im } Q^t$ follows easily from the orthogonality of the last decomposition. \square

5. Strong and weak extensions

Let (M, g) be an oriented Riemannian manifold, let (E_1, h_1) and (E_2, h_2) be Hermitian vector bundles on M , and let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be a differential operator.

Using $\Gamma_0(M, E_j)$ to denote the space of smooth sections of E_j with compact support, we will write the restriction of P to compactly supported sections as $P_0 : \Gamma_0(M, E_1) \rightarrow \Gamma_0(M, E_2)$. This can be viewed as an unbounded, densely defined and closable linear operator $P_0 : L^2 E_1 \rightarrow L^2 E_2$, with domain $\Gamma_0(M, E_1)$. Below, we construct two canonical closed extensions of P_0 , thereby verifying that P_0 is closable.

The *strong extension* P_s (also called the *minimal closed extension* P_{min}) is defined by taking the closure of the graph of P_0 , i.e., $P_s = \overline{P_0}$, or more explicitly

$$\mathcal{D}(P_s) := \{u \in L^2 E_1 \mid \exists \{u_j\}_{j \in \mathbb{N}} \subset \Gamma_0(M, E_1), \exists v \in L^2 E_2 \text{ s.t. } u_j \rightarrow u, P u_j \rightarrow v\},$$

and $P_s u := v$.

The *weak extension* P_w (also called the *maximal closed extension* P_{max}) is defined as the largest extension of P_0 which acts distributionally, i.e.,

$$\mathcal{D}(P_w) := \{u \in L^2 E_1 \mid \exists v \in L^2 E_2, \text{ s.t. } \langle v, w \rangle_2 = \langle u, P^* w \rangle_1, \forall w \in \Gamma_0(M, E_2)\},$$

and $P_w u := v$. Here $\langle \cdot, \cdot \rangle_i$ denotes the inner product defined on $L^2 E_i$. Note that this definition is equivalent to saying P_w is the Hilbert adjoint of P^* restricted to smooth forms with compact support, i.e., $((P^*)_0)^t = P_w$. Moreover $P_s \subseteq P_w$, and every closed extension P' of P_0 which acts distributionally (namely, such that $(P')^t$ is an extension of $(P^*)_0$) is contained between the minimal and the maximal closed extensions.

Remark 5.1. Since a closable densely defined operator and its closure have the same adjoint by Lemma 4.1, it follows that $((P^*)_s)^t = P_w$, thus implying

$$(P^*)_s = (P_w)^t, \quad (P^*)_w = (P_s)^t.$$

Again from Lemma 4.1, we immediately get

$$(\ker P_w)^\perp = \overline{\text{im } P_s^*} \quad (\ker P_s)^\perp = \overline{\text{im } P_w^*}.$$

Moreover, note that if $E_1 = E_2$ and P is formally self-adjoint, then P_0 is essentially self-adjoint if and only if $P_s = P_w$.

Set $E := E_1 = E_2$ and let $P, Q : \Gamma(M, E) \rightarrow \Gamma(M, E)$ be differential operators and denote by $P_0, Q_0 : \Gamma_0(M, E) \rightarrow \Gamma_0(M, E)$ their restrictions to compactly supported sections. It is straightforward to adapt the following theory in the case where P and Q are defined from sections of (E_1, h_1) to sections of another Hermitian vector bundle (E_2, h_2) .

We now provide a number of technical but well-known lemmas which will be needed in subsequent sections.

Lemma 5.2. $\overline{\text{im } P_0} = \overline{\text{im } P_s}$.

Proof. The inclusion \subseteq follows from $\text{im } P_0 \subseteq \text{im } P_s$ passing to the closures. The other inclusion \supseteq follows from $\overline{\text{im } P_0} \supseteq \text{im } P_s$ passing again to the closures. \square

Lemma 5.3. $\overline{\text{im } (QP)_s} \subseteq \overline{\text{im } Q_s}$.

Proof. By Lemma 5.2, $\overline{\text{im } (QP)_s} = \overline{\text{im } (QP)_0}$ and $\overline{\text{im } Q_s} = \overline{\text{im } Q_0}$, therefore it is sufficient to note that $\text{im } (QP)_0 \subseteq \text{im } Q_0$ and then pass to the closures. \square

Lemma 5.4. $\ker P_w \subseteq \ker(QP)_w$.

Proof. If $\alpha \in \ker P_w$, then for every γ smooth section with compact support

$$\langle \alpha, (QP)^* \gamma \rangle = \langle \alpha, P^* Q^* \gamma \rangle = \langle P_w \alpha, Q^* \gamma \rangle = 0,$$

therefore $\alpha \in \ker(QP)_w$. \square

Let us now introduce the order relation $s \leq w$ between strong and weak extensions, as well as the notation $a' = w$ if $a = s$ and $a' = s$ if $a = w$.

Lemma 5.5. If $QP = 0$, then $\overline{\text{im } P_a} \subseteq \ker Q_b$ for any $a, b \in \{s, w\}$ with $a \leq b$.

Proof. It is enough to prove $\text{im } P_a \subseteq \ker Q_a$ and pass to the closures. If $a = s$ this property follows from the definition, while if $a = w$ it follows from the definition and the relation $(QP)^* = P^*Q^*$. \square

Lemma 5.6. *If $QP = 0$, then $\text{im } Q_c^* \perp \text{im } P_a$ for any $a, c \in \{s, w\}$ with $\min(a, c) = s$.*

Proof. The result is equivalent to $\text{im } Q_b^* \perp \text{im } P_a$ for $a, b \in \{s, w\}$ with $a \leq b$, which follows from Theorem 4.2 and Lemma 5.5. \square

Lemma 5.7. *For $a \in \{s, w\}$, we have $P = P_a$ when acting on $\Gamma(M, E) \cap \mathcal{D}(P_a)$.*

Proof. Since $P_s \subseteq P_w$, it is enough to take $\alpha \in \Gamma(M, E) \cap \mathcal{D}(P_w)$ and note that for every $\gamma \in \Gamma_0(M, E)$

$$\langle P\alpha - P_w\alpha, \gamma \rangle = \langle \alpha, P^*\gamma - P^*\gamma \rangle = 0,$$

therefore by the density of $\Gamma_0(M, E)$ in L^2E , the previous equation holds for every γ in L^2 , implying $P\alpha = P_w\alpha$ (for every $x \in M$ it is enough to choose, e.g., $\gamma = (P\alpha - P_w\alpha)\chi_{K_x}$, where χ_{K_x} is the function equal to 1 on the compact $K_x \ni x$ and equal to 0 outside). \square

Lemma 5.8. *It holds that*

$$\mathcal{D}(P_w) \cap \Gamma(M, E) = \{\alpha \in L^2E \cap \Gamma(M, E) \mid P\alpha \in L^2E\}$$

and

$$\ker P_w \cap \Gamma(M, E) = \{\alpha \in L^2E \cap \Gamma(M, E) \mid P\alpha = 0\}.$$

Proof. If $\alpha \in L^2E \cap \Gamma(M, E)$, then for every $\gamma \in \Gamma_0(M, E)$ smooth section with compact support

$$\langle P\alpha, \gamma \rangle = \langle \alpha, P^*\gamma \rangle.$$

Therefore, arguing as in Lemma 5.7, $\alpha \in \mathcal{D}(P_w)$ iff $P\alpha \in L^2E$, and analogously $\alpha \in \ker P_w$ iff $P\alpha = 0$. \square

Remark 5.9. The definitions of strong and weak extensions and the previous technical lemmas can be given in a more general setting.

For any Hilbert space \mathcal{H} that contains $\Gamma_0(M, E)$ as a dense subspace, it is possible to define the strong and weak extensions of $P_0 : \Gamma_0(M, E) \rightarrow \Gamma_0(M, E)$ so long as it admits a formal adjoint with respect to the inner product of \mathcal{H} ; with the same assumptions one can go on to prove Lemmas 5.2-5.6. As a consequence, the L^2 Dolbeault and the L^2

Aeppli-Bott-Chern Hilbert complexes, which will be defined in Sections 6 and 7, may also be constructed for Hilbert spaces other than L^2 . To prove also Lemmas 5.7 and 5.8 in general one would need additional assumptions on \mathcal{H} , so that its inner product looks sufficiently like the L^2 inner product.

In [2] Andreotti and Vesentini introduced a Hilbert space $W^{p,q}$ defined as $\mathcal{D}((\bar{\partial} + \bar{\partial}^*)_s) \cap L^2\Lambda^{p,q}$ on a Hermitian manifold. If the Hermitian metric is Kähler, in [32, Section 5] the second author proved that for a class of differential operators (including $d, d^*, \partial, \partial^*, \bar{\partial}, \bar{\partial}^*$ and all the Laplacians introduced in Section 3) the $W^{p,q}$ formal adjoint coincides with the L^2 formal adjoint. Therefore, Kähler manifolds with the Hilbert space $W^{p,q}$ provide an example of a more general setting where strong and weak extensions can be defined and Lemmas 5.2-5.8 hold.

6. The Dolbeault Hilbert complex

Let (M, g) be a Hermitian manifold of complex dimension n . For any fixed (p, q) , let us consider the following L^2 Dolbeault Hilbert complex

$$\dots \longrightarrow L^2\Lambda^{p,q-1} \xrightarrow{\bar{\partial}_a} L^2\Lambda^{p,q} \xrightarrow{\bar{\partial}_b} L^2\Lambda^{p,q+1} \longrightarrow \dots \tag{14}$$

where $a, b \in \{s, w\}$ with $a \leq b$ and so, by Lemma 5.5, we have $\overline{\text{im } \bar{\partial}_a} \subseteq \ker \bar{\partial}_b$ and consequently the cohomology spaces below are well-defined.

We denote by

$$L^2 H_{\bar{\partial},ab}^{p,q} := \frac{L^2\Lambda^{p,q} \cap \ker \bar{\partial}_b}{L^2\Lambda^{p,q} \cap \overline{\text{im } \bar{\partial}_a}}$$

the associated *unreduced L^2 Dolbeault cohomology*, and by

$$L^2 \bar{H}_{\bar{\partial},ab}^{p,q} := \frac{L^2\Lambda^{p,q} \cap \overline{\ker \bar{\partial}_b}}{L^2\Lambda^{p,q} \cap \overline{\text{im } \bar{\partial}_a}},$$

the associated *reduced L^2 Dolbeault cohomology*.

As described in (11), the Hilbert complex (14) has an associated Laplacian, which is a positive self-adjoint operator given by

$$\Delta_{\bar{\partial},ab} := \bar{\partial}_a \bar{\partial}_a^t + \bar{\partial}_b^t \bar{\partial}_b,$$

defined on $L^2\Lambda^{p,q}$ with domain given by Theorems 4.4 and 4.5. In fact, the operator $\Delta_{\bar{\partial},ab}$ is well-defined even if $a > b$.

Recalling Remark 5.1 and the notation $a' = w$ if $a = s$ and $a' = s$ if $a = w$, note that $\Delta_{\bar{\partial},ab}$ can be rewritten as

$$\Delta_{\bar{\partial},ab} = \bar{\partial}_a \bar{\partial}_{a'}^* + \bar{\partial}_{b'}^* \bar{\partial}_b.$$

With this formulation of $\Delta_{\bar{\partial},ab}$, it is easy to see that it is an extension of the Dolbeault Laplacian $(\Delta_{\bar{\partial}})_0$ acting on the space $A_0^{p,q}$ of smooth compactly supported forms.

By Lemma 4.6 it holds that in $L^2\Lambda^{p,q}$ we have

$$L^2\mathcal{H}_{\bar{\partial},ab}^{p,q} := \ker \Delta_{\bar{\partial},ab} = \ker \bar{\partial}_b \cap \ker \bar{\partial}_a^*, \tag{15}$$

where $L^2\mathcal{H}_{\bar{\partial},ab}^{p,q}$ is the space of L^2 -Dolbeault harmonic forms.

Moreover, by Theorem 4.2, we obtain the following Dolbeault orthogonal decomposition of the Hilbert space $L^2\Lambda^{p,q}$

$$L^2\Lambda^{p,q} = L^2\mathcal{H}_{\bar{\partial},ab}^{p,q} \oplus \overline{\text{im } \bar{\partial}_a} \oplus \overline{\text{im } \bar{\partial}_{b'}^*},$$

$$\ker \bar{\partial}_b = L^2\mathcal{H}_{\bar{\partial},ab}^{p,q} \oplus \overline{\text{im } \bar{\partial}_a}.$$

From this last decomposition we immediately deduce the isomorphism, induced by the identity, between the space of L^2 Dolbeault harmonic forms and L^2 reduced cohomology

$$L^2\mathcal{H}_{\bar{\partial},ab}^{p,q} \simeq L^2\bar{H}_{\bar{\partial},ab}^{p,q}.$$

Note that by elliptic regularity, *i.e.*, Theorem 2.1, we get $L^2\mathcal{H}_{\bar{\partial},ab}^{p,q} \subseteq A^{p,q}$.

Remark 6.1. This theory of the L^2 Dolbeault Hilbert complex holds similarly for the L^2 Hilbert complexes originated by $\partial^2 = 0$ and $d^2 = 0$. In particular, for $a, b \in \{s, w\}$, we can define self-adjoint operators $\Delta_{\partial,ab}$, $\Delta_{d,ab}$, L^2 cohomology spaces $L^2H_{\partial,ab}^{p,q}$, $L^2H_{d,ab}^k$, $L^2\bar{H}_{\partial,ab}^{p,q}$, $L^2\bar{H}_{d,ab}^k$ and spaces of L^2 harmonic forms $L^2\mathcal{H}_{\partial,ab}^{p,q}$, $L^2\mathcal{H}_{d,ab}^k$ with analogous properties.

7. The Aepli-Bott-Chern Hilbert complex

Throughout this section, (M, g) will denote a Hermitian manifold of complex dimension n .

For any fixed bidegree (p, q) , we define the L^2 Aepli-Bott-Chern Hilbert complex, or L^2 ABC complex for short, to be

$$\begin{array}{c}
 \dots \\
 \downarrow \\
 L^2\Lambda^{p-1,q-2} \oplus L^2\Lambda^{p-2,q-1} \\
 \downarrow (\bar{\partial} \oplus \partial)_a \\
 L^2\Lambda^{p-1,q-1} \\
 \downarrow \partial \bar{\partial}_b \\
 L^2\Lambda^{p,q} \\
 \downarrow (\partial + \bar{\partial})_c \\
 L^2\Lambda^{p+1,q} \oplus L^2\Lambda^{p,q+1} \\
 \downarrow \\
 \dots
 \end{array} \tag{16}$$

where $a, b, c \in \{s, w\}$ with $a \leq b \leq c$ denote either strong or weak extensions. Note that we choose to omit the parentheses when writing $\partial \bar{\partial}_b := (\partial \bar{\partial})_b$ for simplicity of notation. The complete definition of the L^2 ABC Hilbert complex is given in Section 10.

By Lemma 5.5 we have

$$\overline{\text{im } (\bar{\partial} \oplus \partial)_a} \subseteq \ker \partial \bar{\partial}_b, \quad \overline{\text{im } \partial \bar{\partial}_b} \subseteq \ker(\partial + \bar{\partial})_c,$$

and therefore the associated *unreduced L^2 Bott-Chern* and *Aeppli cohomology* spaces, defined by

$$L^2 H_{BC,bc}^{p,q} := \frac{L^2 \Lambda^{p,q} \cap \ker(\partial + \bar{\partial})_c}{L^2 \Lambda^{p,q} \cap \overline{\text{im } \partial \bar{\partial}_b}}$$

and

$$L^2 H_{A,ab}^{p-1,q-1} := \frac{L^2 \Lambda^{p-1,q-1} \cap \ker \partial \bar{\partial}_b}{L^2 \Lambda^{p-1,q-1} \cap \overline{\text{im } (\bar{\partial} \oplus \partial)_a}},$$

are well-defined. Similarly, we can define the *reduced L^2 Bott-Chern* and *Aeppli cohomology* spaces

$$L^2 \bar{H}_{BC,bc}^{p,q} := \frac{L^2 \Lambda^{p,q} \cap \ker(\partial + \bar{\partial})_c}{L^2 \Lambda^{p,q} \cap \overline{\text{im } \partial \bar{\partial}_b}}$$

and

$$L^2 \bar{H}_{A,ab}^{p-1,q-1} := \frac{L^2 \Lambda^{p-1,q-1} \cap \ker \partial \bar{\partial}_b}{L^2 \Lambda^{p-1,q-1} \cap \overline{\text{im } (\bar{\partial} \oplus \partial)_a}}.$$

The positive self-adjoint operators associated to the Hilbert complex (16) are

$$\Delta_{BC,bc} := \partial \bar{\partial}_b (\partial \bar{\partial}_b)^t + (\partial + \bar{\partial})_c^t (\partial + \bar{\partial})_c : L^2 \Lambda^{p,q} \rightarrow L^2 \Lambda^{p,q}$$

and

$$\Delta_{A,ab} := (\partial\bar{\partial}_b)^t \partial\bar{\partial}_b + (\bar{\partial} \oplus \partial)_a (\bar{\partial} \oplus \partial)_a^t : L^2\Lambda^{p-1,q-1} \rightarrow L^2\Lambda^{p-1,q-1},$$

with domains given by Theorems 4.4 and 4.5. By analogy with the elliptic complex (4) we also define the following positive self-adjoint operators as L^2 versions of the operators (3).

$$\square_{BC,bc} = \partial\bar{\partial}_b(\partial\bar{\partial}_b)^t + ((\partial + \bar{\partial})_c^t(\partial + \bar{\partial})_c)^2,$$

acting on $L^2\Lambda^{p,q}$, and

$$\square_{A,ab} = (\partial\bar{\partial}_b)^t \partial\bar{\partial}_b + ((\bar{\partial} \oplus \partial)_a(\bar{\partial} \oplus \partial)_a^t)^2,$$

acting on $L^2\Lambda^{p-1,q-1}$. Their domains are given by Theorems 4.4 and 4.5.

Note that without the condition $a \leq b \leq c$, the above operators are still well-defined, however it may not be possible to define the corresponding cohomology spaces.

By Remark 5.1 and recalling the notation $a' = w$ if $a = s$ and $a' = s$ if $a = w$, the operators in the L^2 ABC Hilbert complex have adjoints given by

$$(\bar{\partial} \oplus \partial)_a^t = (\partial^* + \bar{\partial}^*)_{a'}, \quad \partial\bar{\partial}_b^t = \bar{\partial}^* \partial_{b'}^*, \quad (\partial + \bar{\partial})_c^t = (\partial^* \oplus \bar{\partial}^*)_{c'}. \tag{17}$$

This allows us to rewrite the operators $\Delta_{BC,bc}$ and $\Delta_{A,ab}$ as

$$\Delta_{BC,bc} = \partial\bar{\partial}_b \bar{\partial}^* \partial_{b'}^* + (\partial^* \oplus \bar{\partial}^*)_{c'} (\partial + \bar{\partial})_c,$$

$$\Delta_{A,ab} = \bar{\partial}^* \partial_{b'}^* \partial\bar{\partial}_b + (\bar{\partial} \oplus \partial)_a (\partial^* + \bar{\partial}^*)_{a'},$$

and the operators $\square_{BC,bc}$ and $\square_{A,ab}$ as

$$\square_{BC,bc} = \partial\bar{\partial}_b \bar{\partial}^* \partial_{b'}^* + ((\partial^* \oplus \bar{\partial}^*)_{c'} (\partial + \bar{\partial})_c)^2,$$

$$\square_{A,ab} = \bar{\partial}^* \partial_{b'}^* \partial\bar{\partial}_b + ((\bar{\partial} \oplus \partial)_a (\partial^* + \bar{\partial}^*)_{a'})^2.$$

This makes it clear that the above operators are just extensions of the non-elliptic operators $(\Delta_{BC})_0$, $(\Delta_A)_0$ and the elliptic operators $(\square_{BC})_0$, $(\square_A)_0$, acting on the space $A_0^{\bullet,\bullet}$ of smooth forms with compact support.

Remark 7.1. The dual Hilbert complex of (16) is given by

$$\begin{array}{ccc}
 L^2\Lambda^{p-1,q-2} \oplus L^2\Lambda^{p-2,q-1} & & \\
 (\partial^* + \bar{\partial}^*)_{a'} \uparrow & & \\
 L^2\Lambda^{p-1,q-1} & & \\
 \bar{\partial}^* \partial_{b'}^* \uparrow & & (18) \\
 L^2\Lambda^{p,q} & & \\
 (\partial^* + \bar{\partial}^*)_{c'} \uparrow & & \\
 L^2\Lambda^{p+1,q} \oplus L^2\Lambda^{p,q+1} & &
 \end{array}$$

with $a, b, c \in \{s, w\}$ chosen such that $a \leq b \leq c$. This will be called the *dual L^2 ABC Hilbert complex*.

By Lemma 4.6, we see that the kernels of the elliptic and non-elliptic operators coincide, in particular we have

$$\ker \Delta_{BC,bc} = \ker \square_{BC,bc} = \ker \bar{\partial}^* \partial_{b'}^* \cap \ker(\partial + \bar{\partial})_c,$$

on $L^2\Lambda^{p,q}$, and

$$\ker \Delta_{A,ab} = \ker \square_{A,ab} = \ker(\partial^* + \bar{\partial}^*)_{a'} \cap \ker \partial \bar{\partial}_b,$$

on $L^2\Lambda^{p-1,q-1}$. These kernels will be denoted by

$$\begin{aligned}
 L^2\mathcal{H}_{BC,bc}^{p,q} &:= \ker \Delta_{BC,bc} \cap L^2\Lambda^{p,q}, \\
 L^2\mathcal{H}_{A,ab}^{p-1,q-1} &:= \ker \Delta_{A,ab} \cap L^2\Lambda^{p-1,q-1},
 \end{aligned}$$

and we will call them the spaces of L^2 Bott-Chern harmonic forms and L^2 Aeppli harmonic forms, respectively. Since \square_{BC} and \square_A are elliptic and $\ker \Delta_{BC,bc} \subseteq \ker(\square_{BC})_w$, $\ker \Delta_{A,ab} \subseteq \ker(\square_A)_w$, it follows from Theorem 2.1 (elliptic regularity) that these spaces of L^2 harmonic forms are smooth, namely

$$L^2\mathcal{H}_{BC,bc}^{p,q} \subseteq A^{p,q}, \quad L^2\mathcal{H}_{A,ab}^{p,q} \subseteq A^{p,q}.$$

Remark 7.2. As the Hodge $*$ operator is an isometry with respect to the L^2 inner product, the following duality between spaces of harmonic forms is obtained as a consequence of (6).

$$L^2\mathcal{H}_{BC,bc}^{p,q} \simeq L^2\mathcal{H}_{A,c'b'}^{n-q,n-p}.$$

Remark 7.3. The L^2 ABC complex (16) is composed by two separate Hilbert complexes: the L^2 Aeppli Hilbert complex is given by the first two differentials $(\bar{\partial} \oplus \partial)_a$ and $\partial \bar{\partial}_b$, with $a, b \in \{s, w\}$ and $a \leq b$; while the L^2 Bott-Chern Hilbert complex is given by the last two differentials $\partial \bar{\partial}_b$ and $(\partial + \bar{\partial})_c$, with $b, c \in \{s, w\}$ and $b \leq c$.

In the following we might consider these two Hilbert complexes separately. In this way, without any loss of generality, we can increment by $(1, 1)$ the bidegree for the L^2 Aeppli Hilbert complex in (16). As a consequence, we will uniform the subscripts denoting strong or weak extensions in both Hilbert complexes to just a, b with $a \leq b$.

We introduce the notation $L^2 A^{p,q} := L^2 \Lambda^{p,q} \cap A^{p,q}$ to denote the space of L^2 (p, q) -forms which are smooth. The spaces $L^2 A^k$ and $L^2 A^k_{\mathbb{C}}$ can be defined similarly. By applying Theorem 4.2 to the L^2 ABC Hilbert complex, we obtain the following orthogonal decompositions of the Hilbert space $L^2 \Lambda^{p,q}$.

Theorem 7.4 (*L^2 Bott-Chern and Aeppli decompositions*). *The Hilbert space $L^2 \Lambda^{p,q}$ decomposes as*

$$L^2 \Lambda^{p,q} = L^2 \mathcal{H}_{BC,ab}^{p,q} \oplus \overline{\text{im } \partial \bar{\partial}_a} \oplus \overline{\text{im } (\partial^* \oplus \bar{\partial}^*)_{b'}},$$

$$L^2 \Lambda^{p,q} = L^2 \mathcal{H}_{A,ab}^{p,q} \oplus \overline{\text{im } (\bar{\partial} \oplus \partial)_a} \oplus \overline{\text{im } \bar{\partial}^* \partial_{b'}^*},$$

where $a, b \in \{s, w\}$ such that $a \leq b$. We also have

$$\ker(\partial + \bar{\partial})_b = L^2 \mathcal{H}_{BC,ab}^{p,q} \oplus \overline{\text{im } \partial \bar{\partial}_a},$$

$$\ker \partial \bar{\partial}_b = L^2 \mathcal{H}_{A,ab}^{p,q} \oplus \overline{\text{im } (\bar{\partial} \oplus \partial)_a}.$$

Moreover a smooth form $\alpha \in L^2 A^{p,q}$ has smooth components with respect to the above decompositions.

Proof. The decompositions follow from Theorem 4.2, therefore it only remains to prove the regularity. Given $\alpha \in L^2 \Lambda^{p,q}$ we can write

$$\alpha = h + \beta + \eta$$

by the first decomposition, where $h \in \ker \square_{BC,ab}$, $\beta \in \overline{\text{im } \partial \bar{\partial}_a}$ and $\eta \in \overline{\text{im } (\partial^* \oplus \bar{\partial}^*)_{b'}}$. Lemma 5.5 tells us that $\beta \in \ker \partial_w \cap \ker \bar{\partial}_w$ and $\eta \in \ker \bar{\partial}^* \partial_w^*$, which implies that, when α is smooth, β is a weak solution of

$$\square_{BC} \beta = \partial \bar{\partial} \bar{\partial}^* \partial^* \alpha$$

and η is a weak solution of

$$\square_{BC} \eta = (\partial^* \partial + \bar{\partial}^* \bar{\partial})^2 \alpha.$$

By Theorem 2.1, it follows that β and η are also smooth, proving the result. The Aeppli case is analogous. \square

Corollary 7.5. *There exist isomorphisms, induced by the identity, between the spaces of L^2 harmonic forms and reduced L^2 cohomology for $a, b \in \{s, w\}$ such that $a \leq b$*

$$L^2\mathcal{H}_{BC,ab}^{p,q} \simeq L^2\bar{H}_{BC,ab}^{p,q} \qquad L^2\mathcal{H}_{A,ab}^{p,q} \simeq L^2\bar{H}_{A,ab}^{p,q}.$$

By [11, Lemma 3.8] if $(\Delta_{BC})_0$ or $(\Delta_A)_0$ are essentially self-adjoint, then $\partial\bar{\partial}_s = \partial\bar{\partial}_w$ and $\bar{\partial}^*\partial_s^* = \bar{\partial}^*\partial_w^*$. However, since $(\Delta_{BC})_0$ and $(\Delta_A)_0$ are not elliptic, it can be difficult to find actual cases where they are essentially self-adjoint. Instead, we prove that the same result holds when $(\square_{BC})_0$ or $(\square_A)_0$ are essentially self-adjoint; this will become important in Section 9 (see the discussion after Proposition 9.3).

Theorem 7.6. *If $(\square_{BC})_0$ or $(\square_A)_0$ is essentially self-adjoint, then*

$$\partial\bar{\partial}_s = \partial\bar{\partial}_w, \quad \bar{\partial}^*\partial_s^* = \bar{\partial}^*\partial_w^*.$$

Proof. Given a closed linear operator $P : \mathcal{H}_1 \rightarrow \mathcal{H}_2$, a linear subspace $\mathcal{E} \subseteq \mathcal{D}(P)$ is called a *core* if \mathcal{E} is dense in $\mathcal{D}(P)$ with respect to the graph norm $\|x\|_{Gr(P)} := \|x\|_{\mathcal{H}_1} + \|Px\|_{\mathcal{H}_2}$.

By [11, p. 98] we know that

$$\mathcal{E}_{ab} := \bigcap_{k \in \mathbb{N}} \mathcal{D}(\Delta_{A,ab}^k)$$

is a core for $\partial\bar{\partial}_b$ with $a, b \in \{s, w\}$ and $a \leq b$. But by definition $\mathcal{E}_{ab} \subseteq \mathcal{D}(\Delta_{A,ab}^2)$, and by (12) we have $\mathcal{D}(\Delta_{A,ab}^2) \subseteq \mathcal{D}(\square_{A,ab})$, therefore $\mathcal{D}(\square_{A,ab})$ is also a core for $\partial\bar{\partial}_b$. In particular $\mathcal{D}(\square_{A,ss})$ is a core for $\partial\bar{\partial}_s$ and $\mathcal{D}(\square_{A,sw})$ is a core for $\partial\bar{\partial}_w$. Now note that since $\partial\bar{\partial}_s \subseteq \partial\bar{\partial}_w$, we know that $\partial\bar{\partial}_s$ and $\partial\bar{\partial}_w$ coincide on $\mathcal{D}(\square_{A,ss})$, and if $(\square_A)_0$ is essentially self-adjoint then $\mathcal{D}(\square_{A,ss}) = \mathcal{D}(\square_{A,sw})$. By the definition of core this shows $\partial\bar{\partial}_s = \partial\bar{\partial}_w$, and Remark 5.1 ends the proof in the case when $(\square_A)_0$ is essentially self-adjoint.

If $(\square_{BC})_0$ is essentially self-adjoint, considering the dual L^2 ABC Hilbert complex (18), a similar argument shows that $\bar{\partial}^*\partial_s^* = \bar{\partial}^*\partial_w^*$, and again Remark 5.1 ends the proof. \square

Remark 7.7. Actually, [11, Lemma 3.8] yields that if $(\Delta_{BC})_0$ is essentially self-adjoint, then $\partial\bar{\partial}_s = \partial\bar{\partial}_w$ and $(\partial + \bar{\partial})_s = (\partial + \bar{\partial})_w$. Similarly, if $(\Delta_A)_0$ is essentially self-adjoint, then $\partial\bar{\partial}_s = \partial\bar{\partial}_w$ and $(\bar{\partial} \oplus \partial)_s = (\bar{\partial} \oplus \partial)_w$. With the same proof of Theorem 7.6 we can prove that if $(\square_{BC})_0$ is essentially self-adjoint, then $(\partial + \bar{\partial})_s = (\partial + \bar{\partial})_w$, and if $(\square_A)_0$ is essentially self-adjoint, then $(\bar{\partial} \oplus \partial)_s = (\bar{\partial} \oplus \partial)_w$. However, since $\partial + \bar{\partial}$ and $\bar{\partial} \oplus \partial$ are first order differential operators, in actual examples the two properties $(\partial + \bar{\partial})_s = (\partial + \bar{\partial})_w$ and $(\bar{\partial} \oplus \partial)_s = (\bar{\partial} \oplus \partial)_w$ hold when the metric is complete (see, e.g., [1, Theorem 1.3]), which usually is a more frequent assumption to get, compared with essential self-adjointness.

By a similar argument, if $(\square_{BC})_0$ and/or $(\square_A)_0$ are essentially self-adjoint, we are able to characterise the kernel of their unique self-adjoint extension. First note that, as a consequence of Lemma 5.8, we have

$$\begin{aligned} \ker \square_{BC,sw} &= \{\alpha \in L^2 A^{\bullet,\bullet} \mid \partial\alpha = \bar{\partial}\alpha = \bar{\partial}^* \partial^* \alpha = 0\}, \\ \ker \square_{A,sw} &= \{\alpha \in L^2 A^{\bullet,\bullet} \mid \partial^* \alpha = \bar{\partial}^* \alpha = \partial\bar{\partial}\alpha = 0\}. \end{aligned}$$

Furthermore, again by Lemma 5.8, we have

$$\begin{aligned} \ker(\square_{BC})_w &= \{\alpha \in L^2 A^{\bullet,\bullet} \mid \square_{BC}\alpha = 0\}, \\ \ker(\square_A)_w &= \{\alpha \in L^2 A^{\bullet,\bullet} \mid \square_A\alpha = 0\}. \end{aligned}$$

Therefore we obtain the following result.

Proposition 7.8. *If $(\square_{BC})_0$ is essentially self-adjoint, then for any smooth form $\alpha \in L^2 A^{\bullet,\bullet}$ it holds that*

$$\square_{BC}\alpha = 0 \iff \partial\alpha = \bar{\partial}\alpha = \bar{\partial}^* \partial^* \alpha = 0.$$

If $(\square_A)_0$ is essentially self-adjoint, then for any smooth form $\alpha \in L^2 A^{\bullet,\bullet}$ it holds that

$$\square_A\alpha = 0 \iff \partial^* \alpha = \bar{\partial}^* \alpha = \partial\bar{\partial}\alpha = 0.$$

As a direct consequence of [11, Theorem 3.5], the unreduced L^2 Bott-Chern and Aeppli cohomologies can be computed using the subcomplex obtained by intersecting (16) with the space of smooth forms $A^{\bullet,\bullet}$. Note that this subcomplex is not itself a Hilbert complex.

Theorem 7.9 ([11, Theorem 3.5]). *We have the following isomorphisms for $a, b \in \{s, w\}$ and $a \leq b$*

$$\begin{aligned} L^2 H_{BC,ab}^{p,q} &\simeq \frac{\ker(\partial + \bar{\partial})_b \cap A^{p,q}}{\partial\bar{\partial}_a(A^{p-1,q-1} \cap \mathcal{D}(\partial\bar{\partial}_a))}, \\ L^2 H_{A,ab}^{p,q} &\simeq \frac{\ker \partial\bar{\partial}_b \cap A^{p,q}}{(\bar{\partial} \oplus \partial)_a((A^{p,q-1} \oplus A^{p-1,q}) \cap \mathcal{D}((\bar{\partial} \oplus \partial)_a))}. \end{aligned}$$

We note that, in view of Lemma 5.8, when $a = b = c = w$ the statement simplifies to

$$\begin{aligned} L^2 H_{BC,ww}^{p,q} &\simeq \frac{\ker \partial \cap \ker \bar{\partial} \cap L^2 A^{p,q}}{\partial\bar{\partial}(L^2 A^{p-1,q-1}) \cap L^2 A^{p,q}}, \\ L^2 H_{A,ww}^{p,q} &\simeq \frac{\ker \partial\bar{\partial} \cap L^2 A^{p,q}}{(\bar{\partial} \oplus \partial)(L^2 A^{p,q-1} \oplus L^2 A^{p-1,q}) \cap L^2 A^{p,q}}. \end{aligned}$$

As a consequence of the above results, on a compact Hermitian manifold the spaces of Bott-Chern cohomology, Bott-Chern harmonic forms, L^2 reduced Bott-Chern cohomology, L^2 unreduced Bott-Chern cohomology, and L^2 Bott-Chern harmonic forms are all isomorphic.

Indeed, in this setting $(\square_{BC})_0$ is essentially self-adjoint [4, Proposition 4.1] and so $\partial\bar{\partial}_s = \partial\bar{\partial}_w$ by Theorem 7.6. Furthermore $(\partial + \bar{\partial})_s = (\partial + \bar{\partial})_w$ and $(\bar{\partial} \oplus \partial)_s = (\bar{\partial} \oplus \partial)_w$ since the metric g is complete (see, e.g., [1, Theorem 1.3]). Therefore

$$L^2 H_{BC,bc}^{p,q} \simeq L^2 H_{BC,ww}^{p,q}.$$

Moreover, by Theorem 7.9 and Lemma 5.8, we have

$$L^2 H_{BC,ww}^{p,q} \simeq H_{BC}^{p,q},$$

and since M is compact $H_{BC}^{p,q} \simeq \mathcal{H}_{BC}^{p,q}$ has finite dimension. This implies that $\text{im } \partial\bar{\partial}_b$ is closed (cf. [11, Theorem 2.4]), thus

$$L^2 H_{BC,bc}^{p,q} \simeq L^2 \bar{H}_{BC,bc}^{p,q} \simeq L^2 \mathcal{H}_{BC,bc}^{p,q}.$$

Similar isomorphisms also exist for Aeppli cohomology.

7.1. A remark on the reduced L^2 Aeppli and Bott-Chern cohomologies

Recall the smooth Aeppli and Bott-Chern cohomologies, considered in Section 3

$$H_{BC}^{p,q} := \frac{\ker \partial + \bar{\partial}}{\text{im } \partial\bar{\partial}}, \quad H_A^{p,q} := \frac{\ker \partial\bar{\partial}}{\text{im } \bar{\partial} \oplus \partial}.$$

When the bidegree is fixed we can write $\ker \partial + \bar{\partial}$ and $\text{im } \partial\bar{\partial}$ equivalently as $\ker \partial \cap \ker \bar{\partial}$ and $\text{im } \partial + \text{im } \bar{\partial}$. However, things are not so simple for the reduced L^2 Aeppli and Bott-Chern cohomologies

$$L^2 \bar{H}_{BC,ab}^{p,q} := \frac{\ker(\partial + \bar{\partial})_b}{\text{im } \partial\bar{\partial}_a}, \quad L^2 \bar{H}_{A,ab}^{p,q} := \frac{\ker \partial\bar{\partial}_b}{\text{im } (\bar{\partial} \oplus \partial)_a},$$

with $a, b \in \{s, w\}$, $a \leq b$. There is no guarantee that we have $(\partial + \bar{\partial})_b = \partial_b + \bar{\partial}_b$ and as a consequence we cannot assume that $\ker(\partial + \bar{\partial})_b$ and $\text{im } (\bar{\partial} \oplus \partial)_a$ can be used interchangeably with $\ker \partial_b \cap \ker \bar{\partial}_b$ and $\text{im } \bar{\partial}_a + \text{im } \partial_a$; instead we have the following results.

Lemma 7.10. *Let the operators $\partial_a + \bar{\partial}_a$ and $\partial_a^* + \bar{\partial}_a^*$ be defined with domains $\mathcal{D}(\partial_a) \cap \mathcal{D}(\bar{\partial}_a) \subseteq L^2 \Lambda^{p,q}$ and $\mathcal{D}(\partial_a^*) \cap \mathcal{D}(\bar{\partial}_a^*) \subseteq L^2 \Lambda^{p,q}$ respectively, for some choice of strong or weak extensions $a \in \{s, w\}$. Then for the weak extension on $L^2 \Lambda^{p,q}$ we have*

$$(\partial + \bar{\partial})_w = \partial_w + \bar{\partial}_w, \quad (\partial^* + \bar{\partial}^*)_w = \partial_w^* + \bar{\partial}_w^*$$

however for the strong extension we have only

$$(\partial + \bar{\partial})_s \subseteq \partial_s + \bar{\partial}_s, \quad (\partial^* + \bar{\partial}^*)_s \subseteq \partial_s^* + \bar{\partial}_s^*.$$

In particular, on $L^2\Lambda^{p,q}$ we have

$$\ker(\partial + \bar{\partial})_w = \ker \partial_w \cap \ker \bar{\partial}_w, \quad \ker(\partial^* + \bar{\partial}^*)_w = \ker \partial_w^* \cap \ker \bar{\partial}_w^*,$$

for the weak extension but only

$$\ker(\partial + \bar{\partial})_s \subseteq \ker \partial_s \cap \ker \bar{\partial}_s, \quad \ker(\partial^* + \bar{\partial}^*)_s \subseteq \ker \partial_s^* \cap \ker \bar{\partial}_s^*,$$

for the strong extension.

Proof. The key observation is that $\partial_w + \bar{\partial}_w, (\partial + \bar{\partial})_w, \partial_s + \bar{\partial}_s, (\partial + \bar{\partial})_s$ are operators $L^2\Lambda^{p,q} \rightarrow L^2\Lambda^{p+1,q} \oplus L^2\Lambda^{p,q+1}$ and the spaces $L^2\Lambda^{p+1,q}, L^2\Lambda^{p,q+1}$ are orthogonal. Knowing this, both inclusions $\partial_w + \bar{\partial}_w \subseteq (\partial + \bar{\partial})_w$ and $\partial_w + \bar{\partial}_w \supseteq (\partial + \bar{\partial})_w$ follow directly from the definitions, as well as $(\partial + \bar{\partial})_s \subseteq \partial_s \cap \bar{\partial}_s$. The same works also for the operator $\partial^* + \bar{\partial}^*$. \square

An immediate consequence is the following.

Corollary 7.11. *The following equalities hold in $L^2\Lambda^{p,q}$*

$$\overline{\text{im}(\bar{\partial} \oplus \partial)_s} = \overline{\text{im} \partial_s} + \overline{\text{im} \bar{\partial}_s}, \quad \overline{\text{im}(\partial^* \oplus \bar{\partial}^*)_s} = \overline{\text{im} \partial_s^*} + \overline{\text{im} \bar{\partial}_s^*}.$$

Proof. We will just prove the first equality; the second is similar. Lemma 7.10 tells us that we have

$$\ker(\partial^* + \bar{\partial}^*)_w = \ker \partial_w^* \cap \ker \bar{\partial}_w^*,$$

in $L^2\Lambda^{p,q}$. By Remark 5.1, we also have

$$\ker(\partial^* + \bar{\partial}^*)_w = \text{im}(\bar{\partial} \oplus \partial)_s^\perp,$$

and

$$\ker \partial_w^* \cap \ker \bar{\partial}_w^* = \text{im} \partial_s^\perp + \text{im} \bar{\partial}_s^\perp = \left(\overline{\text{im} \partial_s} + \overline{\text{im} \bar{\partial}_s}\right)^\perp.$$

Therefore, the equality

$$\overline{\text{im}(\bar{\partial} \oplus \partial)_s} = \overline{\text{im} \partial_s} + \overline{\text{im} \bar{\partial}_s}$$

holds in $L^2\Lambda^{p,q}$. \square

As a consequence, we have the following equivalent definitions of the reduced L^2 Aeppli and Bott-Chern cohomologies.

Corollary 7.12. *When $a = s$ the reduced L^2 Aeppli cohomology can be written as*

$$L^2 \bar{H}_{A, sb}^{p, q} = \frac{L^2 \Lambda^{p, q} \cap \ker \partial \bar{\partial}_b}{L^2 \Lambda^{p, q} \cap \left(\overline{\text{im } \partial_s + \text{im } \bar{\partial}_s} \right)}.$$

When $b = w$ the reduced L^2 Bott-Chern cohomology can be written as

$$L^2 \bar{H}_{BC, aw}^{p, q} = \frac{L^2 \Lambda^{p, q} \cap \ker \partial_w \cap \ker \bar{\partial}_w}{L^2 \Lambda^{p, q} \cap \overline{\text{im } \partial \bar{\partial}_b}}.$$

This is especially useful when the Hermitian metric g is complete and therefore the strong and the weak extensions of $\partial, \bar{\partial}, \bar{\partial} \oplus \partial$ coincide.

Remark 7.13. Note that, by the same argument as in the proof of Lemma 5.2, we can obtain in $L^2 \Lambda^{p, q}$

$$\overline{\text{im } (\bar{\partial} \oplus \partial)_s} = \overline{\text{im } \partial_s + \text{im } \bar{\partial}_s}, \quad \overline{\text{im } (\partial^* \oplus \bar{\partial}^*)_s} = \overline{\text{im } \partial_s^* + \text{im } \bar{\partial}_s^*},$$

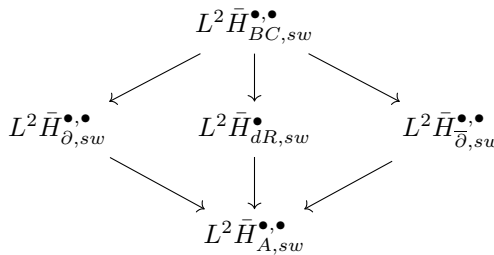
thus Corollary 7.11 implies

$$\overline{\text{im } \partial_s + \text{im } \bar{\partial}_s} = \overline{\text{im } \partial_s + \text{im } \bar{\partial}_s}, \quad \overline{\text{im } \partial_s^* + \text{im } \bar{\partial}_s^*} = \overline{\text{im } \partial_s^* + \text{im } \bar{\partial}_s^*}.$$

7.2. A diagram of maps between the reduced L^2 cohomology spaces

On a complex manifold we have the diagram (5) of maps between complex cohomologies. Given a Hermitian metric, we have a similar diagram involving L^2 reduced cohomology spaces.

Proposition 7.14. *There is the following commutative diagram of maps induced by the identity*



Proof. The maps $L^2 \bar{H}_{BC, sw}^{\bullet, \bullet} \rightarrow L^2 \bar{H}_{\partial, sw}^{\bullet, \bullet}$ and $L^2 \bar{H}_{BC, sw}^{\bullet, \bullet} \rightarrow L^2 \bar{H}_{\bar{\partial}, sw}^{\bullet, \bullet}$ are well-defined since for every bidegree (p, q) the equality

$$\ker(\partial + \bar{\partial})_w = \ker \partial_w \cap \ker \bar{\partial}_w$$

holds in $L^2\Lambda^{p,q}$ by Lemma 7.10, while

$$\overline{\text{im } \partial\bar{\partial}_s} \subseteq \overline{\text{im } \partial_s} \cap \overline{\text{im } \bar{\partial}_s}$$

holds in $L^2\Lambda^{p,q}$ by Lemma 5.3.

The maps $L^2\bar{H}_{\partial,sw}^{\bullet,\bullet} \rightarrow L^2\bar{H}_{A,sw}^{\bullet,\bullet}$ and $L^2\bar{H}_{\bar{\partial},sw}^{\bullet,\bullet} \rightarrow L^2\bar{H}_{A,sw}^{\bullet,\bullet}$ are well-defined since for every bidegree (p, q) the inclusion

$$\ker \partial_w \cup \ker \bar{\partial}_w \subseteq \ker \partial\bar{\partial}_w$$

holds in $L^2\Lambda^{p,q}$ by Lemma 5.4, while $\overline{\text{im } \partial_s} \cup \overline{\text{im } \bar{\partial}_s} \subseteq \overline{\text{im } \partial_s} + \overline{\text{im } \bar{\partial}_s}$ always holds.

The map $L^2\bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2\bar{H}_{dR,sw}^{\bullet,\bullet}$ is well-defined since for every bidegree (p, q) the inclusions

$$\begin{aligned} \ker(\partial + \bar{\partial})_w &\subseteq \ker d_w, \\ \overline{\text{im } \partial\bar{\partial}_s} &\subseteq \overline{\text{im } d_s} \end{aligned}$$

both hold in $L^2\Lambda^{p,q}$ by Lemma 5.3.

Finally the map $L^2\bar{H}_{dR,sw}^{\bullet,\bullet} \rightarrow L^2\bar{H}_{A,sw}^{\bullet,\bullet}$ is well-defined since for every degree k the inclusion

$$\ker d_w \subseteq \bigoplus_{p+q=k} \ker \partial\bar{\partial}_w \cap L^2\Lambda^{p,q}$$

holds in $L^2\Lambda_{\mathbb{C}}^k$ by Lemma 5.4, while

$$\overline{\text{im } d_s} \subseteq \bigoplus_{p+q=k} \overline{\text{im } (\bar{\partial} \oplus \partial)_s} \cap L^2\Lambda^{p,q}$$

holds in $L^2\Lambda_{\mathbb{C}}^k$ by Lemma 5.2. \square

Remark 7.15. The diagram can be generalised replacing the subscripts $_{sw}$ in the first two lines with $_{sb}$, and replacing $_{sw}$ in the last line with $_{aw}$, for $a, b \in \{s, w\}$. Further generalisations would need Lemmas 5.3 and 5.4 to be generalised as well.

Now we observe under which conditions the maps in Proposition 7.14 are isomorphisms.

Proposition 7.16. *The maps in Proposition 7.14 are all isomorphisms if and only if all the following equalities hold in $L^2\Lambda_{\mathbb{C}}^{\bullet}$*

- a) $\overline{\text{im } \partial\bar{\partial}_s} = \ker \partial_w \cap \overline{\ker \bar{\partial}_w} \cap \overline{\text{im } d_s}$,
- b) $\text{im } \partial\bar{\partial}_s = \ker \partial_w \cap \text{im } \bar{\partial}_s$,

- c) $\overline{\text{im } \partial \bar{\partial}_s} = \ker \partial_w \cap \ker \bar{\partial}_w \cap \left(\overline{\text{im } \partial_s} + \overline{\text{im } \bar{\partial}_s} \right),$
- d) $\ker \partial \bar{\partial}_w = \ker d_w + \overline{\text{im } \partial_s} + \overline{\text{im } \bar{\partial}_s},$
- e) $\ker \partial \bar{\partial}_w = \ker \partial_w + \overline{\text{im } \bar{\partial}_s},$
- f) $\ker \partial \bar{\partial}_w = \ker \partial_w \cap \ker \bar{\partial}_w + \overline{\text{im } \partial_s} + \overline{\text{im } \bar{\partial}_s}.$

Note that f) implies e), therefore the above statement still holds if e) is omitted.

Proof. First of all, note that the inclusions \subseteq in conditions a)-f) hold either by Lemma 5.3 or Lemma 5.4. Therefore:

- injectivity of $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{dR,sw}^{\bullet,\bullet}$ is equivalent to a);
- injectivity of $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{\partial,sw}^{\bullet,\bullet}$ is equivalent to b);
- injectivity of $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{\partial,sw}^{\bullet,\bullet}$ is equivalent to the conjugate of b);
- injectivity of $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ is equivalent to c), while surjectivity is equivalent to f);
- surjectivity of $L^2 \bar{H}_{dR,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ is equivalent to d);
- surjectivity of both maps $L^2 \bar{H}_{\partial,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ and $L^2 \bar{H}_{\bar{\partial},sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ follows from f).

To conclude the proof, note that if all the conditions a)-f) hold, then every map starting from $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet}$ is injective and every map arriving in $L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ is surjective. Since the map $L^2 \bar{H}_{BC,sw}^{\bullet,\bullet} \rightarrow L^2 \bar{H}_{A,sw}^{\bullet,\bullet}$ is an isomorphism and the diagram commutes, it follows that every map in the diagram is an isomorphism. \square

The proof of Proposition 7.16 is purely a matter of linear algebra. In fact it is the analogue result of [13, Remark 5.16] involving $\partial, \bar{\partial}, \partial \bar{\partial}$ on smooth forms. Furthermore, in the smooth case the conditions a)-f) are all equivalent to each other [13, Lemma 5.15] (see Section 3). Conversely, here we do not see any simple reason for which the conditions a)-f) in Proposition 7.16 should be equivalent. It would be interesting to know if this is the case or to exhibit counterexamples. We will see in the next section that, on complete Kähler manifolds, all the conditions in Proposition 7.16 hold.

8. Complete Kähler manifolds

Let (M, g) be a complete Kähler manifold. The Kähler assumption guarantees that the relations (7), (8) and (9) all hold. Meanwhile, the completeness of the metric implies that $\delta_s = \delta_w$ for $\delta = d, d^*, \partial, \partial^*, \bar{\partial}, \bar{\partial}^*, \bar{\partial} \oplus \partial, \partial^* \oplus \bar{\partial}^*$; we refer to [1, Theorem 1.3] for a proof of this. Throughout this section, we will write either δ_s or δ_w depending on which is more convenient at the time.

Completeness also implies that all positive integer powers of $(d + d^*)_0, (\partial + \partial^*)_0, (\bar{\partial} + \bar{\partial}^*)_0$ as operators $A_0^{\bullet,\bullet} \rightarrow A_0^{\bullet,\bullet}$ are essentially self-adjoint [12, Section 3.B]. For example, $\Delta_{\bar{\partial},sw}$ is the unique self-adjoint extension of $(\Delta_{\bar{\partial}})_0 = (\bar{\partial} + \bar{\partial}^*)_0^2$. Similarly the operator $\Delta_{\partial,sw}^k$, defined via the Spectral Theorem 4.7, is the unique self-adjoint extension of

$(\Delta_{\bar{\partial}})_0^k = (\bar{\partial} + \bar{\partial}^*)_0^{2k}$. Analogously $\Delta_{\partial,sw}^k$ is the unique self-adjoint extension of $(\Delta_{\partial})_0^k = (\partial + \partial^*)_0^{2k}$ and $\Delta_{d,sw}^k$ is the unique self-adjoint extension of $(\Delta_d)_0^k = (d + d^*)_0^{2k}$. Finally, note that relation (7), together with the essential self-adjointness of the above Laplacians, implies

$$\Delta_{d,sw} = 2\Delta_{\partial,sw} = 2\Delta_{\bar{\partial},sw}. \tag{19}$$

8.1. Equality between L^2 Dolbeault and Bott-Chern harmonic forms

On compact Kähler manifolds, it is known that the kernel of the Dolbeault Laplacian coincides with the kernels of the Bott-Chern and Aeppli Laplacians. We will now show that this statement can be generalised to the complete Kähler case.

We begin by giving a characterisation of the space of L^2 Dolbeault harmonic forms.

Lemma 8.1. *Let (M, g) be a complete Kähler manifold. Then*

$$L^2\mathcal{H}_{\bar{\partial},sw}^{p,q} = \{\alpha \in L^2A^{p,q} \mid \bar{\partial}\alpha = \bar{\partial}^*\alpha = \partial\alpha = \partial^*\alpha = 0\}.$$

Proof. Note that (19) and (15) imply

$$L^2\mathcal{H}_{\bar{\partial},sw}^{p,q} = \ker \bar{\partial}_w \cap \ker \bar{\partial}_w^* \cap \ker \partial_w \cap \ker \partial_w^* \cap L^2\Lambda^{p,q},$$

which is a space of smooth forms by Theorem 2.1. From Lemma 5.8 we conclude the proof. \square

On a general Hermitian manifold (M, g) we define extensions of the Bott-Chern and Aeppli Laplacians $(\tilde{\Delta}_{BC})_0$ and $(\tilde{\Delta}_A)_0$:

$$\begin{aligned} \tilde{\Delta}_{BC,bc} &:= \partial\bar{\partial}_b\bar{\partial}^*\partial_{b'}^* + \bar{\partial}^*\partial_s^*\partial\bar{\partial}_w + \partial^*\bar{\partial}_s\bar{\partial}^*\partial_w + \bar{\partial}^*\partial_s\partial^*\bar{\partial}_w + (\partial^* \oplus \bar{\partial}^*)_{c'}(\partial + \bar{\partial})_c, \\ \tilde{\Delta}_{A,ab} &:= \bar{\partial}^*\partial_{b'}\partial\bar{\partial}_b + \partial\bar{\partial}_s\bar{\partial}^*\partial_w^* + \bar{\partial}\partial_s^*\partial\bar{\partial}_w^* + \partial\bar{\partial}_s^*\bar{\partial}\partial_w^* + (\bar{\partial} \oplus \partial)_a(\partial^* + \bar{\partial}^*)_{a'}, \end{aligned}$$

for $a, b, c \in \{s, w\}$. They are positive and self-adjoint on the domains given by Theorems 4.4 and 4.5. Furthermore, they have the same kernels as the operators $\Delta_{BC,bc}$ and $\Delta_{A,ab}$, respectively.

Lemma 8.2. *Let (M, g) be a Hermitian manifold. For $a, b, c \in \{s, w\}$*

$$\begin{aligned} \ker \tilde{\Delta}_{BC,bc} &= \ker \bar{\partial}^*\partial_{b'}^* \cap \ker(\partial + \bar{\partial})_c = L^2\mathcal{H}_{BC,bc}^{p,q}, \\ \ker \tilde{\Delta}_{A,ab} &= \ker \partial\bar{\partial}_b \cap \ker(\partial^* + \bar{\partial}^*)_{a'} = L^2\mathcal{H}_{A,ab}^{p,q}. \end{aligned}$$

Proof. By Lemma 4.6 we deduce that in $L^2\Lambda^{p,q}$ we have

$$\ker \tilde{\Delta}_{BC,bc} = \ker \bar{\partial}^*\partial_{b'}^* \cap \ker \partial\bar{\partial}_w \cap \ker \bar{\partial}^*\partial_w \cap \ker \partial^*\bar{\partial}_w \cap \ker(\partial + \bar{\partial})_c.$$

The desired result then follows from Lemma 7.10, since

$$\ker(\partial + \bar{\partial})_c \subseteq \ker \partial_c \cap \ker \bar{\partial}_c \subseteq \ker \bar{\partial}_w \cap \ker \partial_w,$$

and from Lemma 5.4, noting that

$$\ker \bar{\partial}_w \subseteq \ker \partial \bar{\partial}_w \cap \ker \partial^* \bar{\partial}_w, \quad \ker \partial_w \subseteq \ker \bar{\partial}^* \partial_w.$$

The Aepli case is analogous. \square

Remark 8.3. If (M, g) is a complete Hermitian manifold, then $(\partial + \bar{\partial})_c = (\partial + \bar{\partial})_w$ and $(\partial^* + \bar{\partial}^*)_{a'} = (\partial^* + \bar{\partial}^*)_w$ by completeness. By Lemma 7.10 we also have $\ker(\partial + \bar{\partial})_w = \ker \partial_w \cap \ker \bar{\partial}_w$ and $\ker(\partial^* + \bar{\partial}^*)_w = \ker \partial_w^* \cap \ker \bar{\partial}_w^*$. Therefore, we can write

$$\ker(\partial + \bar{\partial})_c = \ker \partial_w \cap \ker \bar{\partial}_w, \quad \ker(\partial^* + \bar{\partial}^*)_{a'} = \ker \partial_w^* \cap \ker \bar{\partial}_w^*.$$

Consider the fourth order terms appearing in the operators $\tilde{\Delta}_{BC,bc}$ and $\tilde{\Delta}_{A,ab}$. Taken without the accompanying second order terms, they define the following positive self-adjoint operators:

$$\begin{aligned} \tilde{\Delta}_{BC,b,4} &:= \partial \bar{\partial}_b \bar{\partial}^* \partial_{b'}^* + \bar{\partial}^* \partial_s^* \partial \bar{\partial}_w + \partial^* \bar{\partial}_s \bar{\partial}^* \partial_w + \bar{\partial}^* \partial_s \partial^* \bar{\partial}_w, \\ \tilde{\Delta}_{A,b,4} &:= \bar{\partial}^* \partial_{b'}^* \partial \bar{\partial}_b + \partial \bar{\partial}_s \bar{\partial}^* \partial_w^* + \bar{\partial} \bar{\partial}_s^* \partial \bar{\partial}_w^* + \partial \bar{\partial}_s^* \bar{\partial} \partial_w^*, \end{aligned}$$

with the domains given by Theorems 4.4 and 4.5.

If we assume the Hermitian manifold (M, g) is Kähler and complete then, by (8) and (9), it follows that $\tilde{\Delta}_{BC,b,4}$ and $\tilde{\Delta}_{A,b,4}$ are extensions of $(\Delta_{\bar{\partial}})_0^2$. Since $(\Delta_{\bar{\partial}})_0^2$ is essentially self-adjoint, we then deduce that

$$\tilde{\Delta}_{BC,b,4} = \tilde{\Delta}_{A,b,4} = \Delta_{\bar{\partial},sw}^2.$$

By Lemma 4.6, note that in $L^2 \Lambda^{p,q}$

$$\begin{aligned} \ker \tilde{\Delta}_{BC,b,4} &= \ker \bar{\partial}^* \partial_{b'}^* \cap \ker \partial \bar{\partial}_w \cap \ker \bar{\partial}^* \partial_w \cap \ker \partial^* \bar{\partial}_w, \\ \ker \tilde{\Delta}_{A,b,4} &= \ker \partial \bar{\partial}_b \cap \ker \bar{\partial}^* \partial_w^* \cap \ker \partial \bar{\partial}_w^* \cap \ker \bar{\partial} \partial_w^*, \\ \ker \Delta_{\bar{\partial},sw}^2 &= \ker \Delta_{\bar{\partial},sw}. \end{aligned}$$

Now we can prove that the kernel of the unique self-adjoint extension of the Dolbeault Laplacian coincides with the kernels of our self-adjoint extensions of the Bott-Chern and Aepli Laplacians.

Theorem 8.4. *Let (M, g) be a complete Kähler manifold. Then for $a, b, c \in \{s, w\}$ it holds that*

$$L^2\mathcal{H}_{BC,bc}^{p,q} = L^2\mathcal{H}_{A,ab}^{p,q} = L^2\mathcal{H}_{\bar{\partial},sw}^{p,q}.$$

In particular, the corresponding reduced cohomology spaces are isomorphic

$$L^2\bar{H}_{BC,bc}^{p,q} \simeq L^2\bar{H}_{A,ab}^{p,q} \simeq L^2\bar{H}_{\bar{\partial},sw}^{p,q}.$$

Proof. Let us first prove $\ker \tilde{\Delta}_{BC,bc} \supseteq \ker \Delta_{\bar{\partial},sw}$. By Lemma 8.2 and Remark 8.3, this is the same as proving

$$\ker \bar{\partial}^* \partial_b^* \cap \ker \partial_w \cap \ker \bar{\partial}_w \supseteq \ker \Delta_{\bar{\partial},sw}.$$

By Lemma 8.1 and Lemma 5.8 we get $\ker \partial_w \cap \ker \bar{\partial}_w \supseteq \ker \Delta_{\bar{\partial},sw}$, while $\ker \bar{\partial}^* \partial_b^* \supseteq \ker \Delta_{\bar{\partial},sw}$ follows from $\ker \Delta_{\bar{\partial},sw} = \ker \Delta_{\bar{\partial},sw}^2 = \ker \tilde{\Delta}_{BC,b,4}$.

Conversely, let $\alpha \in \ker \tilde{\Delta}_{BC,bc}$; in particular $\alpha \in \ker \partial_w \cap \ker \bar{\partial}_w$ by Remark 8.3. Then, for every $\gamma \in A_0^{\bullet,\bullet}$, using (8)

$$\begin{aligned} 0 &= \langle \tilde{\Delta}_{BC,bc} \alpha, \gamma \rangle = \langle \alpha, \tilde{\Delta}_{BC} \gamma \rangle, \\ &= \langle \alpha, \Delta_{\bar{\partial}}^2 \gamma + \partial^* \partial \gamma + \bar{\partial}^* \bar{\partial} \gamma \rangle, \\ &= \langle \alpha, \Delta_{\bar{\partial}}^2 \gamma \rangle, \end{aligned}$$

therefore $\alpha \in \ker(\Delta_{\bar{\partial}}^2)_w$. Since $(\Delta_{\bar{\partial}}^2)_0$ is essentially self-adjoint, then $(\Delta_{\bar{\partial}}^2)_w = \Delta_{\bar{\partial},sw}^2$ and so $\ker \Delta_{\bar{\partial},sw} = \ker \Delta_{\bar{\partial},sw}^2$, which allows us to conclude.

In the same way we can prove that $\ker \tilde{\Delta}_{A,ab} = \ker \Delta_{\bar{\partial},sw}$. \square

Corollary 8.5. *If (M, g) is a complete Kähler manifold, then the maps between the L^2 reduced cohomology spaces in Proposition 7.14 are all isomorphisms, and all the conditions a), . . . , f) in Proposition 7.16 hold.*

Proof. By Theorem 8.4 and (19), each class of the reduced L^2 cohomology spaces in Proposition 7.14 have a unique representative in $L^2\mathcal{H}_{d,sw}^\bullet$. Since all the maps are induced by the identity, it follows that they are all isomorphisms. By Proposition 7.16, we get the last claim. \square

An immediate consequence of conditions a), . . . , f) is the following.

Corollary 8.6 (L^2 reduced $\partial\bar{\partial}$ -Lemma on complete Kähler manifolds). *Let (M, g) be a complete Kähler manifold. If $\alpha \in L^2\Lambda_{\mathbb{C}}^\bullet \cap \ker \partial_w \cap \ker \bar{\partial}_w$, then*

$$\alpha \in \overline{\text{im } \partial\bar{\partial}_s} \iff \alpha \in \overline{\text{im } d_s} \iff \alpha \in \overline{\text{im } \partial_s} \iff \alpha \in \overline{\text{im } \bar{\partial}_s} \iff \alpha \in \overline{\text{im } \bar{\partial}_s} + \overline{\text{im } \partial_s}.$$

Corollary 8.7. *Let (M, g) be a complete Kähler manifold. Then*

$$\ker \partial\bar{\partial}_s = \ker \partial\bar{\partial}_w, \quad \overline{\text{im } \partial\bar{\partial}_s} = \overline{\text{im } \partial\bar{\partial}_w},$$

$$\ker \bar{\partial}^* \partial_s^* = \ker \bar{\partial}^* \partial_w^*, \quad \overline{\operatorname{im} \bar{\partial}^* \partial_s^*} = \overline{\operatorname{im} \bar{\partial}^* \partial_w^*}.$$

Proof. By Theorem 7.4, Remark 8.3 and Theorem 8.4 we have

$$\begin{aligned} \ker \partial_w \cap \ker \bar{\partial}_w &= L^2 \mathcal{H}_{\bar{\partial},sw}^{p,q} \oplus \overline{\operatorname{im} \partial \bar{\partial}_a}, \\ \ker \partial \bar{\partial}_b &= L^2 \mathcal{H}_{\bar{\partial},sw}^{p,q} \oplus (\overline{\operatorname{im} \partial_s} + \overline{\operatorname{im} \bar{\partial}_s}). \end{aligned}$$

For different choices of $a, b \in \{s, w\}$, the thesis follows. \square

8.2. Spectral gap of the Dolbeault Laplacian

We will now study the case of a complete Kähler manifold for which $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2 \Lambda^{\bullet,\bullet}$. Specifically, we will explore the consequences this has for L^2 Bott-Chern and Aeppli cohomologies. First of all, note that in this case by Lemma 4.9, $\operatorname{im} \bar{\partial}_s$, $\operatorname{im} \bar{\partial}_s^*$, and $\operatorname{im} \Delta_{\bar{\partial},sw}$ are closed. Second, since $\Delta_{d,sw} = 2\Delta_{\partial,sw} = 2\Delta_{\bar{\partial},sw}$ by (19), it follows that $\operatorname{im} \partial_s$, $\operatorname{im} d_s$, $\operatorname{im} \partial_s^*$ and $\operatorname{im} d_s^*$ are also closed. In particular, if $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2 \Lambda^{\bullet,\bullet}$, then the reduced and unreduced L^2 cohomologies coincide for both the Dolbeault and de Rham cases, *i.e.*, for $a, b \in \{s, w\}$ with $a \leq b$, we have

$$L^2 \bar{H}_{\bar{\partial},ab}^{\bullet,\bullet} = L^2 H_{\bar{\partial},ab}^{\bullet,\bullet}, \quad L^2 \bar{H}_{\partial,ab}^{\bullet,\bullet} = L^2 H_{\partial,ab}^{\bullet,\bullet}, \quad L^2 \bar{H}_{d,ab}^{\bullet} = L^2 H_{d,ab}^{\bullet}.$$

We want to conclude that the same holds for L^2 Bott-Chern and Aeppli cohomology. In proving this, we will also prove that a spectral gap of $\Delta_{\bar{\partial},sw}$ implies a spectral gap of $\Delta_{A,ab}$ and $\Delta_{BC,ab}$. A simple proof of this result does not seem viable and so, over the course of the next couple pages, we will work our way towards it, primarily through the use of Lemma 4.9.

We start by proving that $\operatorname{im} (\bar{\partial} \oplus \partial)_s$ and $\operatorname{im} (\partial^* \oplus \bar{\partial}^*)_s$ are closed in this setting.

Theorem 8.8. *Let (M, g) be a complete Kähler manifold. If $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2 \Lambda^{p,q-1} \oplus L^2 \Lambda^{p-1,q}$, then $\operatorname{im} (\bar{\partial} \oplus \partial)_s$ is closed in $L^2 \Lambda^{p,q}$. If $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2 \Lambda^{p+1,q} \oplus L^2 \Lambda^{p,q+1}$, then $\operatorname{im} (\partial^* \oplus \bar{\partial}^*)_s$ is closed in $L^2 \Lambda^{p,q}$.*

Proof. We just prove the first claim since the second one is analogous. Our strategy is to use Lemma 4.9, therefore we consider the following Hilbert complex:

$$L^2 \Lambda^{p,q-2} \oplus L^2 \Lambda^{p-1,q-1} \oplus L^2 \Lambda^{p-2,q} \xrightarrow{(\bar{\partial} \oplus d \oplus \partial)_s} L^2 \Lambda^{p,q-1} \oplus L^2 \Lambda^{p-1,q} \xrightarrow{(\bar{\partial} \oplus \partial)_w} L^2 \Lambda^{p,q}.$$

We consider an operator $\square : A^{p,q-1} \oplus A^{p-1,q} \rightarrow A^{p,q-1} \oplus A^{p-1,q}$ defined by

$$\square := (\partial^* + \bar{\partial}^*)(\bar{\partial} \oplus \partial) + (\bar{\partial} \oplus d \oplus \partial)(\partial^* + \bar{\partial}^*),$$

where $\bar{\partial} \oplus d \oplus \partial$ acts on $A^{p,q-2} \oplus A^{p-1,q-1} \oplus A^{p-2,q}$ component by component. Define also $\square_{sw} : L^2 \Lambda^{p,q-1} \oplus L^2 \Lambda^{p-1,q} \rightarrow L^2 \Lambda^{p,q-1} \oplus L^2 \Lambda^{p-1,q}$ by

$$\square_{sw} := (\partial^* + \bar{\partial}^*)_s(\bar{\partial} \oplus \partial)_w + (\bar{\partial} \oplus d \oplus \partial)_s(\partial^* + \bar{\partial}^*)_w,$$

so that \square_{sw} is the Laplacian associated to the above Hilbert complex. By Theorems 4.4 and 4.5 it is a self-adjoint extension of the operator \square_0 on $A_0^{p,q-1} \oplus A_0^{p-1,q}$. By Lemma 4.9, to prove the theorem it is sufficient to show that \square_{sw} has a spectral gap. Note that by the Kähler identities, namely using $\partial\bar{\partial}^* + \bar{\partial}^*\partial = 0$ and $\bar{\partial}\partial^* + \partial^*\bar{\partial} = 0$, we have

$$\square(\alpha + \beta) = \Delta_{\bar{\partial}}(\alpha + \beta) + \partial\partial^*\alpha + \bar{\partial}\bar{\partial}^*\beta$$

for all $\alpha \in A^{p,q-1}$ and $\beta \in A^{p-1,q}$. Arguing as in [41, Theorem 2.4], we can show that \square_0 is essentially self-adjoint. Therefore, if by Theorems 4.4 and 4.5 we define the self-adjoint operator

$$\square'_{sw} := \Delta_{\bar{\partial},sw} + \partial_s\partial_w^* \oplus \bar{\partial}_s\bar{\partial}_w^* : L^2\Lambda^{p,q-1} \oplus L^2\Lambda^{p-1,q} \rightarrow L^2\Lambda^{p,q-1} \oplus L^2\Lambda^{p-1,q},$$

then $\square_{sw} = \square'_{sw}$. We are left to prove that \square'_{sw} has a spectral gap. Note that in $L^2\Lambda^{p,q-1} \oplus L^2\Lambda^{p-1,q}$ it holds that $\mathcal{D}(\square'_{sw}) = \mathcal{D}(\Delta_{\bar{\partial},sw})$, and $\ker \square'_{sw} = \ker \Delta_{\bar{\partial},sw}$ by Lemma 8.1. Then there exists a $C > 0$ such that for all $(\alpha \oplus \beta) \in \mathcal{D}(\square'_{sw}) \cap (\ker \square'_{sw})^\perp$

$$\begin{aligned} C\|\alpha + \beta\|^2 &\leq \langle \Delta_{\bar{\partial},sw}(\alpha + \beta), \alpha + \beta \rangle \\ &\leq \langle \Delta_{\bar{\partial},sw}(\alpha + \beta), \alpha + \beta \rangle + \|\partial_w^*\alpha\|^2 + \|\bar{\partial}_w^*\beta\|^2 \\ &= \langle \Delta_{\bar{\partial},sw}(\alpha + \beta), \alpha + \beta \rangle + \langle \partial_s\partial_w^*\alpha, \alpha \rangle + \langle \bar{\partial}_s\bar{\partial}_w^*\beta, \beta \rangle \\ &= \langle \Delta_{\bar{\partial},sw}(\alpha + \beta) + \partial_s\partial_w^*\alpha + \bar{\partial}_s\bar{\partial}_w^*\beta, \alpha + \beta \rangle \\ &= \langle \square'_{sw}(\alpha + \beta), \alpha + \beta \rangle, \end{aligned}$$

namely \square'_{sw} has a spectral gap. \square

Remark 8.9. The Hilbert complex considered in the previous proof is nothing but the preceding stage of the L^2 ABC Hilbert complex as shown in (16) (see Section 10 for more details).

Now we deal with the closedness of $\text{im } \partial\bar{\partial}_s$ and $\text{im } \bar{\partial}^*\partial_s^*$. By the Spectral Theorem 4.7, we see that if $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2\Lambda^{p,q}$, then $\Delta_{\bar{\partial},sw}^k$ also has a spectral gap in $L^2\Lambda^{p,q}$ for all $k \in \mathbb{N}$, $k \geq 2$. Namely, if $C = \inf(\sigma(\Delta_{\bar{\partial},sw}) \setminus \{0\}) > 0$, then $C^k = \inf(\sigma(\Delta_{\bar{\partial},sw}^k) \setminus \{0\}) > 0$. In particular, for $b \in \{s, w\}$, since $\Delta_{\bar{\partial},sw}^2 = \tilde{\Delta}_{BC,b,4} = \tilde{\Delta}_{A,b,4}$, then $\tilde{\Delta}_{BC,b,4}$ and $\tilde{\Delta}_{A,b,4}$ both also have a spectral gap. The next theorem, the proof of which is inspired by [21, Theorem 1.4.A], shows that $\text{im } \partial\bar{\partial}_s$ and $\text{im } \bar{\partial}^*\partial_s^*$ are closed in this setting.

Theorem 8.10. *Let (M, g) be a complete Kähler manifold. If $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2\Lambda^{p,q}$, then $\text{im } \partial\bar{\partial}_b$ is closed in $L^2\Lambda^{p+1,q+1}$ and $\text{im } \bar{\partial}^*\partial_b^*$ is closed in $L^2\Lambda^{p-1,q-1}$ for $b \in \{s, w\}$.*

Proof. By the previous discussion and Lemma 4.8, since $\tilde{\Delta}_{A,w,4}$ has a spectral gap, there is a constant $C > 0$ such that

$$\|\partial\bar{\partial}_w\psi\|^2 + \|\bar{\partial}^*\partial_w^*\psi\|^2 + \|\partial\bar{\partial}_w^*\psi\|^2 + \|\bar{\partial}\partial_w^*\psi\|^2 = \langle\psi, \tilde{\Delta}_{A,w,4}\psi\rangle \geq C\langle\psi, \psi\rangle \tag{20}$$

for all $\psi \in \mathcal{D}(\tilde{\Delta}_{A,w,4}) \cap (\ker \tilde{\Delta}_{A,w,4})^\perp \cap L^2\Lambda^{p,q}$.

Recall that $\text{im } \partial\bar{\partial}_b = \text{im } \partial\bar{\partial}_s$ by Corollary 8.7. Let $\alpha \in \overline{\text{im } \partial\bar{\partial}_s} \cap L^2\Lambda^{p+1,q+1} = \overline{\partial\bar{\partial}A_0^{p,q}}$, then there is a sequence of $\beta_k \in A_0^{p,q}$ such that $\partial\bar{\partial}\beta_k \rightarrow \alpha$. By the Aepli decompositions of Theorem 7.4, we obtain

$$L^2A^{p,q} = \ker \partial\bar{\partial}_w \cap A^{p,q} \oplus \overline{\text{im } \bar{\partial}^*\partial_s^*} \cap A^{p,q},$$

therefore we can decompose

$$\beta_k = \eta_k + \theta_k,$$

where $\eta_k \in \ker \partial\bar{\partial}_w \cap A^{p,q}$ and $\theta_k \in \overline{\text{im } \bar{\partial}^*\partial_s^*} \cap A^{p,q}$. In particular $\theta_k \in (\ker \tilde{\Delta}_{A,w,4})^\perp$ and $\partial\bar{\partial}\beta_k = \partial\bar{\partial}\eta_k + \partial\bar{\partial}\theta_k = \partial\bar{\partial}\theta_k$ by Lemma 5.8. By Lemma 5.5, we also know that $\theta_k \in \ker \bar{\partial}_w^* \cap \ker \partial_w^*$, therefore $\bar{\partial}^*\theta_k = \partial^*\theta_k = 0$ by Lemma 5.8, implying that $\theta_k \in \mathcal{D}(\tilde{\Delta}_{A,sw,4})$.

Applying (20) we find that

$$C\|\theta_k - \theta_j\|^2 \leq \|\partial\bar{\partial}\theta_k - \partial\bar{\partial}\theta_j\|^2 \rightarrow 0$$

as $k, j \rightarrow +\infty$ since the sequence $\partial\bar{\partial}\beta_k = \partial\bar{\partial}\theta_k$ is Cauchy. From this we conclude that the sequence θ_k is also Cauchy and so it must converge to some form $\theta \in L^2\Lambda^{p,q}$, proving that $\partial\bar{\partial}_s\theta = \alpha$. In particular $\partial\bar{\partial}_b\theta = \alpha$.

The proof for $\text{im } \bar{\partial}^*\partial_b^*$ is analogous, using the Bott-Chern decompositions of Theorem 7.4 and the spectral gap property for $\tilde{\Delta}_{BC,s,4}$. \square

Corollary 8.11. *Let (M, g) be a complete Kähler manifold. If $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2\Lambda^{\bullet,\bullet}$, then for $a, b \in \{s, w\}$ with $a \leq b$*

$$L^2\bar{H}_{BC,ab}^{\bullet,\bullet} = L^2H_{BC,ab}^{\bullet,\bullet}, \quad L^2\bar{H}_{A,ab}^{\bullet,\bullet} = L^2H_{A,ab}^{\bullet,\bullet}.$$

Applying Lemma 4.9, we finally see that a spectral gap of $\Delta_{\bar{\partial},sw}$ actually implies a spectral gap also of $\Delta_{A,ab}$ and $\Delta_{BC,ab}$.

Corollary 8.12. *Let (M, g) be a complete Kähler manifold. If $\Delta_{\bar{\partial},sw}$ has a spectral gap in $L^2\Lambda^{\bullet,\bullet}$, then $\Delta_{A,ab}$ and $\Delta_{BC,ab}$ have a spectral gap in $L^2\Lambda^{\bullet,\bullet}$, with $a, b \in \{s, w\}$ and $a \leq b$.*

Corollary 8.13. *Let (M, g) be a complete Kähler manifold. If $\Delta_{\bar{\partial},s_w}$ has a spectral gap in $L^2\Lambda^{\bullet,\bullet}$, then*

$$\partial\bar{\partial}_s = \partial\bar{\partial}_w, \quad \bar{\partial}^* \partial_s^* = \bar{\partial}^* \partial_w^*.$$

Proof. By [6, Proposition 1.6], if $\ker \partial\bar{\partial}_s = \ker \partial\bar{\partial}_w$ and $\text{im } \partial\bar{\partial}_s = \text{im } \partial\bar{\partial}_w$, then $\partial\bar{\partial}_s = \partial\bar{\partial}_w$. This follows from Corollary 8.7 and Theorem 8.10. The same argument shows that we also have $\bar{\partial}^* \partial_s^* = \bar{\partial}^* \partial_w^*$. \square

In particular, if on a complete Kähler manifold $\Delta_{\bar{\partial},s_w}$ has a spectral gap in $L^2\Lambda^{\bullet,\bullet}$, then there is a unique L^2 ABC Hilbert complex.

We end the section recalling a sufficient condition for $\Delta_{\bar{\partial},s_w}$ to have a spectral gap, due to Gromov.

Theorem 8.14 ([21, Theorems 1.2.B, 1.4.A]). *Let (M, g) be a complete Kähler manifold of complex dimension n . If the fundamental form ω is d -bounded, i.e., $\omega = d\eta$ with η bounded, then $\ker \Delta_{\bar{\partial},s_w} \cap L^2\Lambda^{p,q} = \{0\}$ if $p + q \neq n$ and $\Delta_{\bar{\partial},s_w}$ has a spectral gap in $L^2\Lambda^{\bullet,\bullet}$.*

9. Galois coverings of a compact complex manifold

Given a smooth manifold M , we say that two Riemannian metrics $g^{(1)}$ and $g^{(2)}$ on M are *quasi-isometric* if there exists a positive real constant C such that at every point $p \in M$ and for every tangent vector $X_p \in T_pM$ we have

$$\frac{1}{C}g_p^{(2)}(X_p, X_p) \leq g_p^{(1)}(X_p, X_p) \leq Cg_p^{(2)}(X_p, X_p). \tag{21}$$

Now let M be a complex manifold. As outlined in Section 3, any Hermitian metric on M induces a Hermitian metric on the bundle of (p, q) -forms $\Lambda^{p,q}M$. Furthermore, if we have two quasi-isometric Hermitian metrics $g^{(1)}$ and $g^{(2)}$ on M , the induced metrics $h^{(1)}$ and $h^{(2)}$ will satisfy the same inequality, i.e., there exists a constant such that at any point $p \in M$ a general (p, q) -form $\alpha_p \in \Lambda_p^{p,q}M$ satisfies

$$\frac{1}{C}h_p^{(2)}(\alpha_p, \alpha_p) \leq h_p^{(1)}(\alpha_p, \alpha_p) \leq Ch_p^{(2)}(\alpha_p, \alpha_p).$$

This pointwise inequality implies a more general inequality for the L^2 norm. Specifically, for some constant $C > 0$, any (not necessarily smooth) (p, q) -form $\alpha : M \rightarrow \Lambda^{p,q}M$ satisfies

$$\frac{1}{C}\|\alpha\|_{L^2,h^{(2)}} \leq \|\alpha\|_{L^2,h^{(1)}} \leq C\|\alpha\|_{L^2,h^{(2)}}.$$

From this we see that, although quasi-isometry may change the norm of $L^2\Lambda^{p,q}$, the underlying set of L^2 (p, q) -forms does not change. Furthermore, quasi-isometry preserves

the notion of convergence: if a sequence within $L^2\Lambda^{p,q}$ converges for one Hermitian metric, it converges for all metrics in the same quasi-isometry class. It follows that if P is a differential operator that doesn't depend on the metric (e.g., $d, \partial, \bar{\partial}, \partial\bar{\partial}$) the strong extension P_s defined with respect to two quasi-isometric metrics is the same operator. It is well known that this same result holds also for the weak extension P_w ; it follows from a metric-independent generalisation of the notion of the formal adjoint, which will not be treated here. However, we will see later in this section that, in the setting we are interested in, $P_s = P_w$ for $P \in \{d, \partial, \bar{\partial}, \partial\bar{\partial}\}$, therefore the previous discussion for P_s is sufficient to show that also P_w depends only on the quasi-isometry class of the metric.

From this we can obtain the following result for Bott-Chern and Aeppli cohomology.

Proposition 9.1. *On a Hermitian manifold, the reduced and unreduced L^2 Bott-Chern cohomologies, $L^2\bar{H}_{BC,bc}^{p,q}$ and $L^2H_{BC,bc}^{p,q}$, and also the reduced and unreduced L^2 Aeppli cohomologies, $L^2\bar{H}_{A,ab}^{p,q}$ and $L^2H_{A,ab}^{p,q}$, depend only on the quasi-isometry class of the Hermitian metric and the choice of closed extensions $a, b, c \in \{s, w\}$ with $a \leq b \leq c$ in the L^2 Aeppli-Bott-Chern complex.*

By the same argument it is well known that the above proposition holds for other L^2 cohomologies defined on differential forms, such as the de Rham or Dolbeault cohomology.

Following from the discussion after Theorem 7.9 we know that, on a compact complex manifold, the L^2 Aeppli and Bott-Chern cohomologies are uniquely defined and coincide with the cohomologies defined on smooth forms. In this section, we consider the more general case of a Galois covering of a compact complex manifold.

Let $\pi : \widetilde{M} \rightarrow M$ be a covering of a smooth manifold M with Γ denoting the group of deck transformations. If Γ acts transitively on the fibre $\pi^{-1}(p)$ for all points $p \in M$ and \widetilde{M} is connected, then we say that π is a *Galois Γ -covering* and \widetilde{M} is a *Galois Γ -covering space*. One example of a Galois covering is provided by the universal covering.

Conversely, if Γ is a discrete group acting freely and properly discontinuously on a smooth connected manifold \widetilde{M} , then \widetilde{M}/Γ is a smooth manifold and the map

$$\pi : \widetilde{M} \rightarrow \widetilde{M}/\Gamma$$

is a Galois Γ -covering.

The main theorem of Galois theory on coverings states the following. Let M be a connected manifold. Given a connected covering $\widetilde{M} \rightarrow M$, we obtain a subgroup $\pi_1(\widetilde{M}) \subset \pi_1(M)$. This defines a bijection between the isomorphism classes of connected coverings of M and the subgroups $G \subset \pi_1(M)$. Under this correspondence, a Galois Γ -covering corresponds to the normal subgroup $G \subset \pi_1(M)$ for which Γ is isomorphic to $\pi_1(M)/G$, and \widetilde{M} is diffeomorphic to the quotient of the universal covering of M by G .

Suppose now that M is a complex manifold. By taking the pullback with respect to the covering map π of the complex structure of M , a Galois Γ -covering space \widetilde{M} inherits

a complex structure. We will now show that, when M is a compact complex manifold, the L^2 Aeppli and Bott-Chern cohomology spaces can be uniquely defined on \widetilde{M} with the induced complex structure. By Proposition 9.1 this means we must rule out dependence on the Hermitian metric and on the choice of closed extensions in the Hilbert complex.

We shall first consider any dependence the cohomology spaces have on the Hermitian metric. In particular, we will show that the L^2 cohomology spaces on \widetilde{M} are the same for any choice of Γ -invariant metric on \widetilde{M} . A metric on \widetilde{M} is said to be Γ -invariant if it is invariant under the pullback with respect to any $\gamma \in \Gamma$, namely if $g = \gamma^*g$, i.e., if at any point $p \in \widetilde{M}$ we have

$$g_p(X_p, Y_p) = g_{\gamma(p)}(d_p\gamma(X_p), d_p\gamma(Y_p))$$

for all $\gamma \in \Gamma$ and all $X_p, Y_p \in T_p\widetilde{M}$. Note that a metric \tilde{g} is Γ -invariant if and only if it is given by the pullback of a metric g on M , that is $\tilde{g} = \pi^*g$.

We prove the following proposition.

Proposition 9.2. *If M is compact, then any two Γ -invariant Riemannian metrics on a Galois Γ -covering space \widetilde{M} are quasi-isometric.*

Proof. We first show that any two metrics on a compact manifold are quasi-isometric.

Let $g^{(1)}, g^{(2)}$ be a pair of Riemannian metrics. We can define the unit sphere bundle $SM \subset TM$ to be the set of unit vectors with respect to the first metric,

$$SM = \left\{ X_p \in T_pM \mid g_p^{(1)}(X_p, X_p) = 1, p \in M \right\}.$$

The second metric then defines a function on SM , $X_p \mapsto g_p^{(2)}(X_p, X_p)$. Since SM is a compact manifold, this function has a well-defined maximum and minimum

$$A := \min_{X_p \in SM} g_p^{(2)}(X_p, X_p) \qquad B := \max_{X_p \in SM} g_p^{(2)}(X_p, X_p).$$

Moreover, since $g_p^{(2)}(X_p, X_p)$ is positive for all X_p , A and B are also positive.

Any element of TM can then be given by re-scaling an element of SM , and thus if we choose $C = \max\{\frac{1}{A}, B\}$ we have the inequality

$$\frac{1}{C}g_p^{(1)}(X_p, X_p) \leq g_p^{(2)}(X_p, X_p) \leq Cg_p^{(1)}(X_p, X_p)$$

for all $p \in M$ and all $X_p \in T_pM$.

The pullback of $g^{(1)}$ and $g^{(2)}$ to \widetilde{M} defines a pair of Γ -invariant metrics given by

$$(\pi^*g^{(i)})_{\tilde{p}}(X_{\tilde{p}}, Y_{\tilde{p}}) = g_{\pi(\tilde{p})}^{(i)}(d_{\tilde{p}}\pi(X_{\tilde{p}}), d_{\tilde{p}}\pi(Y_{\tilde{p}}))$$

for $i = 1, 2$ and for all $\tilde{p} \in \widetilde{M}$ and $X_{\tilde{p}}, Y_{\tilde{p}} \in T_{\tilde{p}}\widetilde{M}$. Since $g^{(1)}$ and $g^{(2)}$ are quasi-isometric, it follows directly from the definition that $\pi^*g^{(1)}$ and $\pi^*g^{(2)}$ are also quasi-isometric. \square

Now, we focus on the closed extensions of the Hilbert complex. Any metric on a compact manifold is complete. Taking the pullback onto the Γ -covering space \widetilde{M} we see likewise that any Γ -invariant metric on \widetilde{M} is complete.

From this we can conclude that the strong and weak extensions coincide for first order differential operators such as $\overline{\partial} \oplus \partial$ and $\partial + \overline{\partial}$ in the ABC complex. See, e.g., [1, Theorem 1.3] for a proof. Unfortunately, this result cannot be applied to second order differential operators such as $\partial \overline{\partial}$. Instead we will make use of the next Proposition, which can be applied to general elliptic operators on vector bundles, combined with Theorem 7.6.

Let $P : \Gamma(M, E_1) \rightarrow \Gamma(M, E_2)$ be a differential operator between two vector bundles E_1 and E_2 on a smooth manifold M . Given a Galois Γ -covering, we can lift P to a differential operator \widetilde{P} between smooth sections of the induced pullback bundles \widetilde{E}_1 and \widetilde{E}_2 on the Γ -covering space \widetilde{M} as follows. Any $\alpha \in \Gamma(\widetilde{M}, \widetilde{E}_1)$ can be locally written as $\pi^* \beta$ with $\beta \in \Gamma(M, E_1)$, and this allows us to define the global operator \widetilde{P} by its local action $\widetilde{P}\pi^* \beta := \pi^* P\beta$. Basically the expression of P in local coordinates gives a local formula for the lift \widetilde{P} . Note that the differential operator \widetilde{P} is Γ -equivariant, in the sense that

$$\gamma^*(\widetilde{P}\alpha) = \widetilde{P}(\gamma^*\alpha)$$

for all $\gamma \in \Gamma$ and $\alpha \in \Gamma(\widetilde{M}, \widetilde{E}_1)$.

Proposition 9.3 ([5], Proposition 3.1). *Let (E, h) be a Hermitian vector bundle on a smooth manifold M and let $\pi : \widetilde{M} \rightarrow M$ be a Galois Γ -covering. Given an elliptic operator $P : \Gamma(M, E) \rightarrow \Gamma(M, E)$, denote its lift to \widetilde{M} by \widetilde{P} . If M is compact then $\widetilde{P}_s = \widetilde{P}_w$.*

Given a Hermitian metric on M , let the Galois Γ -covering space \widetilde{M} be endowed with the Γ -invariant pullback metric. Since $\pi^* d = d\pi^*$ and both the complex structure and the Hermitian metric on \widetilde{M} are the pullback of the respective structures on M , it follows that $\pi^* \square_{BC} = \square_{BC} \pi^*$. In particular the elliptic operator \square_{BC} on \widetilde{M} is just the lift of the same operator on M . By the above proposition we see that $(\square_{BC})_s = (\square_{BC})_w$ on \widetilde{M} , i.e., using Remark 5.1, $(\square_{BC})_0$ is essentially self-adjoint on \widetilde{M} . Theorem 7.6 then tells us that the strong and weak extensions of $\partial \overline{\partial}$ coincide, namely $\partial \overline{\partial}_s = \partial \overline{\partial}_w$.

From the above discussion we can now conclude the following.

Theorem 9.4. *Let $\pi : \widetilde{M} \rightarrow M$ be a Galois Γ -covering with M a compact complex manifold. There exists a unique L^2 Aeppli-Bott-Chern complex on \widetilde{M} , which is determined by the induced complex structure and by any Γ -invariant metric on \widetilde{M} .*

The corresponding unreduced and reduced, Aeppli and Bott-Chern cohomologies, which we will denote respectively by

$$L^2 H_{A,\Gamma}^{\bullet,\bullet}(M), \quad L^2 H_{BC,\Gamma}^{\bullet,\bullet}(M) \quad \text{and} \quad L^2 \bar{H}_{A,\Gamma}^{\bullet,\bullet}(M), \quad L^2 \bar{H}_{BC,\Gamma}^{\bullet,\bullet}(M)$$

are similarly unique and depend only on the Galois Γ -covering.

Given any Hermitian vector bundle (E, h) on M , we denote the Hermitian pullback bundle on a Galois Γ -covering space \widetilde{M} by $(\widetilde{E}, \widetilde{h})$, with $\widetilde{E} = \pi^*E$ and $\widetilde{h} = \pi^*h$. Note that a Hermitian metric \widetilde{h} on \widetilde{E} is the pullback of a Hermitian metric h on E iff \widetilde{h} is Γ -invariant, i.e., if $\widetilde{h} = \gamma^*\widetilde{h}$ for all $\gamma \in \Gamma$, that is if

$$\widetilde{h}_{\tilde{x}}(\alpha(\tilde{x}), \alpha(\tilde{x})) = \widetilde{h}_{\gamma(\tilde{x})}(((\gamma^{-1})^*\alpha)(\gamma(\tilde{x})), ((\gamma^{-1})^*\alpha)(\gamma(\tilde{x})))$$

for all $\tilde{x} \in \widetilde{M}$, α a section of \widetilde{E} and $\gamma \in \Gamma$.

In the following we consider the space of square-integrable sections $L^2\widetilde{E}$ with respect to any Γ -invariant metric (note that any two Γ -invariant metrics are quasi-isometric). The group Γ then acts as isometries on $L^2\widetilde{E}$ given by the pullback $\gamma^*\alpha$ for any $\gamma \in \Gamma$, $\alpha \in L^2\widetilde{E}$.

A closed Hilbert subspace $V \subseteq L^2\widetilde{E}$ is called a *Hilbert Γ -module* if it is preserved by the action of Γ , i.e., $\gamma^*V = V$ for any $\gamma \in \Gamma$. Maps between Hilbert Γ -modules are given by bounded Γ -equivariant operators, namely bounded operators P which commute with the pullback with respect to any $\gamma \in \Gamma$, i.e., $\gamma^*(P\alpha) = P(\gamma^*\alpha)$ for any $\alpha \in L^2\widetilde{E}$. A map between Hilbert Γ -modules is called a *weak isomorphism* if it is injective and has dense image. A *strong isomorphism* is moreover an isometric isomorphism. It is well known, see, e.g., [17, Lemma 2.5.3], that if there is a weak isomorphism between Hilbert Γ -modules then a strong isomorphism can be built by the polar decomposition. A sequence of maps between Hilbert Γ -modules

$$\dots \longrightarrow V_{i-1} \xrightarrow{d_{i-1}} V_i \xrightarrow{d_i} V_{i+1} \longrightarrow \dots$$

is called *weakly exact* if $\overline{\text{im } d_{i-1}} = \ker d_i$ for all i .

We will now define the Γ -dimension of a Hilbert Γ -module $V \subseteq L^2\widetilde{E}$. We refer to [5,17,28] for a detailed treatment.

First take an orthonormal basis $\{\varphi_i\}$ of V and define the function on \widetilde{M} ,

$$\tilde{f}(\tilde{x}) = \sum_i \tilde{h}_{\tilde{x}}(\varphi_i(\tilde{x}), \varphi_i(\tilde{x})).$$

Note that the choice of orthonormal basis $\{\varphi_i\}$ in the definition is not important. Indeed, given a different basis $\{\psi_j\}$ we can write

$$\begin{aligned} \tilde{f}(\tilde{x}) &= \sum_i \tilde{h}_{\tilde{x}}(\varphi_i(\tilde{x}), \varphi_i(\tilde{x})) \\ &= \sum_i \sum_j \tilde{h}_{\tilde{x}}(\varphi_i(\tilde{x}), \psi_j(\tilde{x})) \\ &= \sum_j \tilde{h}_{\tilde{x}}(\psi_j(\tilde{x}), \psi_j(\tilde{x})). \end{aligned}$$

Swapping the order of summation is allowed here as all of the summands are positive. For any $\gamma \in \Gamma$ we have

$$\gamma^* \tilde{f}(\tilde{x}) = \tilde{f}(\gamma(\tilde{x})) = \sum_i \tilde{h}_{\gamma(\tilde{x})}(\varphi_i(\gamma(\tilde{x})), \varphi_i(\gamma(\tilde{x}))),$$

and since V is a Hilbert Γ -module then $\{(\gamma^{-1})^* \varphi_i\}$ provides an orthonormal basis of V , so that

$$\gamma^* \tilde{f}(\tilde{x}) = \sum_i \tilde{h}_{\gamma(\tilde{x})}(((\gamma^{-1})^* \varphi_i)(\gamma(\tilde{x})), ((\gamma^{-1})^* \varphi_i)(\gamma(\tilde{x}))) = \tilde{f}(\tilde{x}),$$

using that \tilde{h} is Γ -invariant. Therefore \tilde{f} descends to a well-defined function f on M . The Γ -dimension is then given by

$$\dim_{\Gamma} V = \int_M f(x) d\mu.$$

Note that in the case when Γ is the trivial group, the Γ -dimension simply counts the number of elements in the basis, *i.e.*, it is the usual dimension of V as a vector space. Another immediate property of the Γ -dimension is that $\dim_{\Gamma} V = 0$ if and only if $V = \{0\}$.

In principle, the value of \tilde{f} may be infinite at some or all points in \tilde{M} , in which case the Γ -dimension may also be infinite. However by the following proposition, we see that in certain cases we can guarantee that \tilde{f} is finite at all points, and thus when M is compact, the Γ -dimension is finite.

Proposition 9.5 ([5, Proposition 2.4]). *Given a smooth manifold \tilde{M} along with a Hermitian vector bundle (\tilde{E}, \tilde{h}) and an elliptic differential operator acting on smooth sections $\tilde{P} : \Gamma(\tilde{M}, \tilde{E}) \rightarrow \Gamma(\tilde{M}, \tilde{E})$, consider the kernel of the weak extension, $\ker \tilde{P}_w \subseteq L^2 \tilde{E}$. Given any orthonormal basis $\{\varphi_i\}$ for $\ker \tilde{P}_w$, the series*

$$\tilde{f}(\tilde{x}) = \sum_i \tilde{h}(\varphi_i(\tilde{x}), \varphi_i(\tilde{x})),$$

and all its derivatives, converge uniformly on compact subsets of \tilde{M} . Consequently, \tilde{f} is a smooth function.

For any Γ -invariant Hermitian metric on \tilde{M} , Theorem 7.4 implies the isomorphisms

$$\begin{aligned} L^2 \bar{H}_{A,\Gamma}^{p,q}(M) &\cong L^2 \mathcal{H}_{A,sw}^{p,q}, \\ L^2 \bar{H}_{BC,\Gamma}^{p,q}(M) &\cong L^2 \mathcal{H}_{BC,sw}^{p,q}. \end{aligned}$$

Recall the definition of $L^2\mathcal{H}_{A,sw}^{p,q}$ and $L^2\mathcal{H}_{sw}^{p,q}$ as the kernels of $\square_{A,sw}$ and $\square_{BC,sw}$. Proposition 9.3 tells us that $(\square_A)_0$ and $(\square_{BC})_0$ are essentially self-adjoint operators. The choice of closed extensions $\square_{A,sw}$, $\square_{BC,sw}$ is therefore arbitrary but will be convenient later on.

Since the differential operators $\partial, \bar{\partial}, \partial\bar{\partial}$ on \widetilde{M} and their formal adjoints are Γ -equivariant, and the Hermitian metric on $\Lambda^{p,q}\widetilde{M}$ is Γ -invariant, it is easy to check that $L^2\mathcal{H}_{BC,sw}^{p,q}$ and $L^2\mathcal{H}_{A,sw}^{p,q}$ are Hilbert Γ -modules.

We can therefore define the *L² Aeppli* and *Bott-Chern numbers*, denoted by $h_{A,\Gamma}^{p,q}$ and $h_{BC,\Gamma}^{p,q}$, to be

$$h_{A,\Gamma}^{p,q}(M) = \dim_{\Gamma} L^2\mathcal{H}_{A,sw}^{p,q}, \quad h_{BC,\Gamma}^{p,q}(M) = \dim_{\Gamma} L^2\mathcal{H}_{BC,sw}^{p,q}.$$

By Proposition 9.5 the Aeppli and Bott-Chern numbers are finite whenever M is compact.

The space of Aeppli harmonic L^2 forms $L^2\mathcal{H}_{A,sw}^{p,q}$ in general depends on the choice of Γ -invariant Hermitian metric. When necessary we will write $L^2\mathcal{H}_{A,sw}^{p,q}(g)$ to denote the space corresponding to the metric g , and likewise for the space of Bott-Chern harmonic L^2 forms. Regardless of this, as a consequence of the next proposition, the L^2 Aeppli and Bott-Chern numbers are still independent of the Hermitian metric on the compact complex manifold M . We will need the following fundamental property of the Γ -dimension, which will be also used later on.

Lemma 9.6 ([17, Corollary 3.4.6]). *Let*

$$0 \longrightarrow V_1 \longrightarrow V_2 \longrightarrow \dots \longrightarrow V_n \longrightarrow 0$$

be a weakly exact sequence of Hilbert Γ -modules. Then

$$\sum_k (-1)^k \dim_{\Gamma} V_k = 0$$

Proposition 9.7. *Given any two Γ -invariant Hermitian metrics g_1 and g_2 on \widetilde{M} , we have*

$$\begin{aligned} \dim_{\Gamma} L^2\mathcal{H}_{A,sw}^{p,q}(g_1) &= \dim_{\Gamma} L^2\mathcal{H}_{A,sw}^{p,q}(g_2), \\ \dim_{\Gamma} L^2\mathcal{H}_{BC,sw}^{p,q}(g_1) &= \dim_{\Gamma} L^2\mathcal{H}_{BC,sw}^{p,q}(g_2). \end{aligned}$$

Proof. There exists an isomorphism $j : L^2\mathcal{H}_{BC,sw}^{p,q}(g_1) \rightarrow L^2\mathcal{H}_{BC,sw}^{p,q}(g_2)$, constructed via the isomorphism with the reduced cohomology

$$L^2\mathcal{H}_{BC,sw}^{p,q}(g_i) \simeq L^2\bar{H}_{BC,\Gamma}^{p,q}$$

for $i = 1, 2$. In fact we now show that the isomorphism j is actually realised by the projection

$$\pi_{L^2\mathcal{H}_{BC,sw}^{p,q}(g_2)} : L^2\mathcal{H}_{BC,sw}^{p,q}(g_1) \rightarrow L^2\mathcal{H}_{BC,sw}^{p,q}(g_2).$$

First we remark that $\overline{\text{im } \partial\bar{\partial}_s}(g_1) = \overline{\text{im } \partial\bar{\partial}_s}(g_2)$ since g_1 and g_2 are quasi isometric, therefore we can write $\overline{\text{im } \partial\bar{\partial}_s}$ without any confusion. Any form $\alpha_1 \in L^2\mathcal{H}_{BC,sw}^{p,q}(g_1)$ can be decomposed, by Theorem 7.4, as $\alpha_1 = \alpha_2 + \beta$, with $\alpha_2 \in L^2\mathcal{H}_{BC,sw}^{p,q}(g_2)$ and $\beta \in \overline{\text{im } \partial\bar{\partial}_s}$. Therefore j first sends α_1 in the class $[\alpha_1]_{L^2\bar{H}_{BC,\Gamma}^{p,q}}$ which in turn is sent to α_2 . Thus $j(\alpha_1) = \alpha_2$.

Our goal is to show that j is a weak isomorphism of Hilbert Γ -modules, so that then by applying the previous Lemma we end the proof. Being j an isomorphism, it is clearly injective and it has dense image. Moreover j is a bounded operator since g_1 and g_2 are quasi isometric.

Finally, to see that j is Γ -equivariant, we must show that $j(\gamma^*\alpha_1) = \gamma^*(j(\alpha_1))$ for all $\alpha_1 \in L^2\mathcal{H}_{BC,sw}^{p,q}(g_1)$. Since \square_{BC} is Γ -equivariant, then $\gamma^*\alpha_2 \in L^2\mathcal{H}_{BC,sw}^{p,q}(g_2)$. Therefore our claim is equivalent to showing that $\gamma^*\beta \in \overline{\text{im } \partial\bar{\partial}_s}$. This is true since by Lemma 5.2 the form β is the L^2 limit of the smooth and compactly supported forms $\partial\bar{\partial}\beta_j$ and the forms $\partial\bar{\partial}\gamma^*\beta_j$ are as well smooth and compactly supported satisfying

$$\|\gamma^*\beta - \partial\bar{\partial}\gamma^*\beta_j\| = \|\gamma^*\beta - \gamma^*\partial\bar{\partial}\beta_j\| = \|\beta - \partial\bar{\partial}\beta_j\| \rightarrow 0.$$

A similar weak isomorphism of Hilbert Γ -modules exists the same way between $L^2\mathcal{H}_{A,sw}^{p,q}(g_i)$ for $i = 1, 2$. \square

Remark 9.8. By following the same reasoning as above but replacing the L^2 Aeppli-Bott-Chern Hilbert complex with the L^2 de Rham or Dolbeault Hilbert complex, it is possible to define the unique L^2 reduced cohomology spaces $L^2\bar{H}_{d,\Gamma}^k(M)$, $L^2\bar{H}_{\partial,\Gamma}^{p,q}(M)$ and $L^2\bar{H}_{\bar{\partial},\Gamma}^{p,q}(M)$, from which we obtain the L^2 Betti numbers $b_\Gamma^k(M)$, and the L^2 Hodge numbers, $h_{\partial,\Gamma}^{p,q}(M)$ and $h_{\bar{\partial},\Gamma}^{p,q}(M)$.

Given a Galois Γ -covering space \widetilde{M} of a complex manifold M , we define the following subspaces of $L^2\Lambda^{\bullet,\bullet}$

$$\begin{aligned} \mathcal{A}^{\bullet,\bullet} &= \overline{\text{im } \bar{\partial}_s} \cap \overline{\text{im } \partial_s} \cap \ker \bar{\partial}_w^* \partial_w^*, & \mathcal{B}^{\bullet,\bullet} &= \ker \bar{\partial}_w \cap \overline{\text{im } \bar{\partial}_s} \cap \ker \bar{\partial}_s^* \partial_w^*, \\ \mathcal{C}^{\bullet,\bullet} &= \ker \partial\bar{\partial}_w \cap \overline{\text{im } \bar{\partial}_s^*} \cap \ker \partial_w^*, & \mathcal{D}^{\bullet,\bullet} &= \overline{\text{im } \bar{\partial}_s} \cap \ker \partial_w \cap \ker \bar{\partial}_s^* \partial_w^*, \\ \mathcal{E}^{\bullet,\bullet} &= \ker \partial\bar{\partial}_w \cap \overline{\text{im } \bar{\partial}_s^*} \cap \ker \bar{\partial}_w^*, & \mathcal{F}^{\bullet,\bullet} &= \ker \partial\bar{\partial}_w \cap \overline{\text{im } \bar{\partial}_s^*} \cap \overline{\text{im } \bar{\partial}_s^*}. \end{aligned}$$

These spaces actually are Hilbert Γ -modules. It is easy to check it using Lemma 5.2 and the next facts: the differential operators $\partial, \bar{\partial}, \partial\bar{\partial}$ on \widetilde{M} and their formal adjoints are Γ -equivariant; the Hermitian metric on $\Lambda^{p,q}\widetilde{M}$ is Γ -invariant.

Remark 9.9. Note that there are isomorphisms induced by the identity between the spaces just defined and the following quotient spaces in $L^2\Lambda^{\bullet,\bullet}$:

$$\begin{aligned}
 \mathcal{A}^{\bullet,\bullet} \simeq A^{\bullet,\bullet} &:= \frac{\overline{\text{im } \bar{\partial}_s} \cap \overline{\text{im } \partial_s}}{\text{im } \partial \bar{\partial}_s}, & \mathcal{B}^{\bullet,\bullet} \simeq B^{\bullet,\bullet} &:= \frac{\overline{\text{im } \bar{\partial}_s} \cap \ker \bar{\partial}_w}{\text{im } \partial \bar{\partial}_s}, \\
 \mathcal{C}^{\bullet,\bullet} \simeq C^{\bullet,\bullet} &:= \frac{\ker \partial \bar{\partial}_w}{\ker \bar{\partial}_w + \overline{\text{im } \bar{\partial}_s}}, & \mathcal{D}^{\bullet,\bullet} \simeq D^{\bullet,\bullet} &:= \frac{\overline{\text{im } \bar{\partial}_s} \cap \ker \partial_w}{\text{im } \partial \bar{\partial}_s}, \\
 \mathcal{E}^{\bullet,\bullet} \simeq E^{\bullet,\bullet} &:= \frac{\ker \partial \bar{\partial}_w}{\ker \partial_w + \overline{\text{im } \bar{\partial}_s}}, & \mathcal{F}^{\bullet,\bullet} \simeq F^{\bullet,\bullet} &:= \frac{\ker \partial \bar{\partial}_w}{\ker \partial_w + \ker \bar{\partial}_w}.
 \end{aligned}$$

The following sequences are defined just by compositions of natural inclusions and projections (see the proof of the next proposition for the actual definitions of the maps in (22); the maps in (23) are defined similarly)

$$0 \longrightarrow \mathcal{A}^{\bullet,\bullet} \longrightarrow \mathcal{B}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{A,sw}^{\bullet,\bullet} \longrightarrow \mathcal{C}^{\bullet,\bullet} \longrightarrow 0 \tag{22}$$

$$0 \longrightarrow \mathcal{D}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{BC,sw}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} \longrightarrow \mathcal{E}^{\bullet,\bullet} \longrightarrow \mathcal{F}^{\bullet,\bullet} \longrightarrow 0 \tag{23}$$

which actually are maps between Hilbert Γ -modules: it can be shown arguing similarly to the last part of the proof of Proposition 9.7.

We prove that the two sequences above are exact.

Proposition 9.10. *The sequences (22) and (23) are exact sequences of Hilbert Γ -modules.*

Proof. We just prove the exactness of the sequence (22) and omit the proof for (23) since it is very similar. We divide the proof in three steps.

Step 1): we build the maps of the sequence as a combination of projections and inclusions.

$\mathcal{A}^{\bullet,\bullet} \longrightarrow \mathcal{B}^{\bullet,\bullet}$: This map is simply an inclusion.

$\mathcal{B}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet}$: This map is given by composing the inclusion

$$B^{\bullet,\bullet} \rightarrow \ker \bar{\partial}_w$$

with the projection

$$\ker \bar{\partial}_w \rightarrow L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} = \ker \bar{\partial}_w \cap \ker \bar{\partial}_w^*.$$

$L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} \longrightarrow L^2\mathcal{H}_{A,sw}^{\bullet,\bullet}$: This map is given by composing the inclusion

$$L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} \rightarrow \ker \partial \bar{\partial}_w$$

with the projection

$$\ker \partial \bar{\partial}_w \rightarrow L^2\mathcal{H}_{A,sw}^{\bullet,\bullet} = \ker \partial_w^* \cap \ker \bar{\partial}_w^* \cap \ker \partial \bar{\partial}_w.$$

$L^2\mathcal{H}_{A,sw}^{\bullet,\bullet} \rightarrow \mathcal{C}^{\bullet,\bullet}$: This map is just a projection.

Step 2): there is a commutative diagram between sequences

$$\begin{array}{ccccccccc}
 0 & \rightarrow & \mathcal{A}^{\bullet,\bullet} & \rightarrow & \mathcal{B}^{\bullet,\bullet} & \rightarrow & L^2\mathcal{H}_{\bar{\partial},sw}^{\bullet,\bullet} & \rightarrow & L^2\mathcal{H}_{A,sw}^{\bullet,\bullet} & \rightarrow & \mathcal{C}^{\bullet,\bullet} & \rightarrow & 0 \\
 & & \downarrow \simeq & & \downarrow \simeq & & \downarrow \simeq & & \downarrow \simeq & & \downarrow \simeq & & \\
 0 & \rightarrow & A^{\bullet,\bullet} & \rightarrow & B^{\bullet,\bullet} & \rightarrow & L^2\bar{H}_{\bar{\partial},sw}^{\bullet,\bullet} & \rightarrow & L^2\bar{H}_{A,sw}^{\bullet,\bullet} & \rightarrow & C^{\bullet,\bullet} & \rightarrow & 0
 \end{array}$$

where the maps of the sequence in the bottom line are just induced by the identity.

Step 3): the sequence in the bottom line is exact. It is straightforward to verify it once one uses the equivalent definition of $L^2\bar{H}_{A,sw}^{\bullet,\bullet}$ given by Corollary 7.12. \square

We refer to [43, Section 3] for the original idea behind the previous exact sequences. In [3, Theorem A] Angella and Tomassini, starting from the exact sequences in [43, Section 3], were able to prove an inequality on compact manifolds between the Aeppli, Bott-Chern and Dolbeault numbers,

$$h_{\bar{\partial}}^{p,q} + h_{\bar{\partial}}^{p,q} \leq h_A^{p,q} + h_{BC}^{p,q}.$$

We will conclude this section by proving the same inequality holds for the L^2 Aeppli, Bott-Chern and Dolbeault numbers on any Galois Γ -covering of a compact complex manifold. However we must first make a few comments about the exact sequences (22) and (23).

We will write the Γ -dimension of the spaces $\mathcal{A}^{p,q}, \mathcal{B}^{p,q}, \mathcal{C}^{p,q}, \mathcal{D}^{p,q}, \mathcal{E}^{p,q}$ and $\mathcal{F}^{p,q}$ as $a^{p,q}, b^{p,q}, c^{p,q}, d^{p,q}, e^{p,q}$ and $f^{p,q}$. We can see that $\mathcal{A}^{p,q}, \mathcal{B}^{p,q}$ and $\mathcal{D}^{p,q}$ are all contained within $L^2\mathcal{H}_{BC,sw}^{p,q}$, while $\mathcal{C}^{p,q}, \mathcal{E}^{p,q}$ and $\mathcal{F}^{p,q}$ are contained within $L^2\mathcal{H}_{A,sw}^{p,q}$. Thus, by a simple property of the Γ -dimension given in the proposition below, we have $a^{p,q}, b^{p,q}, d^{p,q} \leq h_{A,\Gamma}^{p,q}(M)$ and $c^{p,q}, e^{p,q}, f^{p,q} \leq h_{BC,\Gamma}^{p,q}(M)$. In particular, if M is compact, then $a^{p,q}, b^{p,q}, c^{p,q}, d^{p,q}, e^{p,q}$ and $f^{p,q}$ are all finite.

Proposition 9.11. *If $U, V \subset L^2E$ are Hilbert Γ -modules such that $U \subset V$, then $\dim_{\Gamma} U \leq \dim_{\Gamma} V$.*

Proof. An orthonormal basis $\{\varphi_i\}_{i \in I}$ of V can be chosen such that $\{\varphi_j\}_{j \in J}$ is an orthonormal basis of U , for some subset $J \subset I$. Clearly we have

$$\sum_{j \in J} h_{\tilde{x}}(\varphi_j(\tilde{x}), \varphi_j(\tilde{x})) \leq \sum_{i \in I} h_{\tilde{x}}(\varphi_i(\tilde{x}), \varphi_i(\tilde{x}))$$

for all \tilde{x} and so it follows from the definition that $\dim_{\Gamma} U \leq \dim_{\Gamma} V$. \square

The following is the main result of this section.

Theorem 9.12. For any Galois Γ -covering $\pi : \widetilde{M} \rightarrow M$ of a compact complex manifold M the L^2 Aeppli, Bott-Chern and Dolbeault numbers satisfy the inequality

$$h_{\overline{\partial},\Gamma}^{p,q}(M) + h_{\partial,\Gamma}^{p,q}(M) \leq h_{A,\Gamma}^{p,q}(M) + h_{BC,\Gamma}^{p,q}(M).$$

Proof. We will follow a similar argument to the proof of Theorem A in [3].

First note that by complex conjugation we obtain the equalities

$$\begin{aligned} a^{p,q} &= a^{q,p}, & b^{p,q} &= d^{q,p}, & c^{p,q} &= e^{q,p}, & f^{p,q} &= f^{q,p}, \\ h_{A,\Gamma}^{p,q} &= h_{A,\Gamma}^{q,p}, & h_{BC,\Gamma}^{p,q} &= h_{BC,\Gamma}^{q,p}, & h_{\overline{\partial},\Gamma}^{p,q} &= h_{\overline{\partial},\Gamma}^{q,p}. \end{aligned}$$

From this, Proposition 9.10 and Lemma 9.6, we find that

$$\begin{aligned} h_{BC,\Gamma}^{p,q} + h_{A,\Gamma}^{p,q} &= h_{BC,\Gamma}^{p,q} + h_{A,\Gamma}^{q,p} \\ &= h_{\overline{\partial},\Gamma}^{p,q} + h_{\overline{\partial},\Gamma}^{q,p} + a^{q,p} - b^{q,p} + c^{q,p} + d^{p,q} - e^{p,q} + f^{p,q} \\ &= h_{\overline{\partial},\Gamma}^{p,q} + h_{\overline{\partial},\Gamma}^{p,q} + a^{p,q} + f^{p,q} \\ &\geq h_{\overline{\partial},\Gamma}^{p,q} + h_{\overline{\partial},\Gamma}^{p,q}, \end{aligned}$$

which is the desired result. \square

Considering that we have actually proved $h_{BC,\Gamma}^{p,q} + h_{A,\Gamma}^{p,q} = h_{\overline{\partial},\Gamma}^{p,q} + h_{\overline{\partial},\Gamma}^{p,q} + a^{p,q} + f^{p,q}$, and therefore in particular $\sum_{p+q=k} h_{\overline{\partial},\Gamma}^{p,q} + h_{\overline{\partial},\Gamma}^{p,q} = \sum_{p+q=k} h_{A,\Gamma}^{p,q} + h_{BC,\Gamma}^{p,q} + a^{p,q} + f^{p,q}$, we immediately obtain the following two corollaries.

Corollary 9.13. For any Galois Γ -covering $\pi : \widetilde{M} \rightarrow M$ of a compact complex manifold M it holds that

$$h_{\overline{\partial},\Gamma}^{p,q}(M) + h_{\partial,\Gamma}^{p,q}(M) = h_{A,\Gamma}^{p,q}(M) + h_{BC,\Gamma}^{p,q}(M)$$

if and only if in $L^2\Lambda^{p,q}\widetilde{M}$

$$\ker \partial\overline{\partial}_w = \ker \partial_w + \ker \overline{\partial}_w, \quad \overline{\text{im } \partial\overline{\partial}_s} = \overline{\text{im } \partial_s} \cap \overline{\text{im } \overline{\partial}_s}.$$

Corollary 9.14. For any Galois Γ -covering $\pi : \widetilde{M} \rightarrow M$ of a compact complex manifold M it holds that

$$\sum_{p+q=k} h_{\overline{\partial},\Gamma}^{p,q}(M) + h_{\partial,\Gamma}^{p,q}(M) = \sum_{p+q=k} h_{A,\Gamma}^{p,q}(M) + h_{BC,\Gamma}^{p,q}(M)$$

if and only if in $L^2\Lambda_{\mathbb{C}}^k\widetilde{M}$

$$\ker \partial\overline{\partial}_w = \ker \partial_w + \ker \overline{\partial}_w, \quad \overline{\text{im } \partial\overline{\partial}_s} = \overline{\text{im } \partial_s} \cap \overline{\text{im } \overline{\partial}_s}.$$

Finally, if the compact complex manifold M admits a Kähler metric, the corresponding pullback metric on \widetilde{M} is complete and Kähler, therefore by Theorem 8.4 we deduce the following.

Corollary 9.15. *For any Galois Γ -covering $\pi : \widetilde{M} \rightarrow M$ of a compact Kähler manifold M for any (p, q) it holds that*

$$h_{\overline{\partial}, \Gamma}^{p,q}(M) + h_{\partial, \Gamma}^{p,q}(M) = h_{A, \Gamma}^{p,q}(M) + h_{BC, \Gamma}^{p,q}(M).$$

10. Further remarks and open questions

Let us present the full definition of the ABC complex and the L^2 ABC complex, which we have delayed until now since it was not really necessary for our purposes.

Given a complex manifold of dimension n and a couple of integers (p, q) such that $0 \leq p, q \leq n$, we define the spaces

$$\begin{aligned} \mathcal{L}_{p,q}^k &:= \bigoplus_{r+s=k, r < p, s < q} A^{r,s} && \text{for } k \leq p + q - 2, \\ \mathcal{L}_{p,q}^{k-1} &:= \bigoplus_{r+s=k, r \geq p, s \geq q} A^{r,s} && \text{for } k \leq p + q, \end{aligned}$$

as well as the differentials $\delta_{p,q}^k : \mathcal{L}_{p,q}^k \rightarrow \mathcal{L}_{p,q}^{k+1}$ with

$$\begin{aligned} \delta_{p,q}^k &= \pi_{\mathcal{L}_{p,q}^{k+1}} \circ d && \text{for } k \leq p + q - 3, \\ \delta_{p,q}^k &= \partial \overline{\partial} && \text{for } k = p + q - 2, \\ \delta_{p,q}^k &= d && \text{for } k \geq p + q - 1, \end{aligned}$$

where $\pi_{\mathcal{L}_{p,q}^{k+1}}$ is the projection onto $\mathcal{L}_{p,q}^{k+1}$. The full ABC complex is then

$$0 \longrightarrow \mathcal{L}_{p,q}^0 \xrightarrow{\delta_{p,q}^0} \mathcal{L}_{p,q}^1 \xrightarrow{\delta_{p,q}^1} \dots \xrightarrow{\delta_{p,q}^{2n-2}} \mathcal{L}_{p,q}^{2n-1} \xrightarrow{\delta_{p,q}^{2n-1}} \mathcal{L}_{p,q}^{2n} \longrightarrow 0.$$

We have already observed in Section 3 that this complex is elliptic. In particular, given a Hermitian metric g on M , the Laplacians $\Delta_{p,q}^k : \mathcal{L}_{p,q}^k \rightarrow \mathcal{L}_{p,q}^k$ defined by (3) are elliptic. If M is compact, it follows that

$$\mathcal{H}_{p,q}^k := \ker \Delta_{p,q}^k = \ker \delta_{p,q}^k \cap \ker (\delta_{p,q}^{k-1})^*,$$

where $(\delta_{p,q}^{k-1})^*$ denotes the formal adjoint of $\delta_{p,q}^{k-1}$, and all the dimensions

$$h_{p,q}^k := \dim_{\mathbb{C}} \mathcal{H}_{p,q}^k$$

are finite (cf. [40]).

We can similarly define the spaces of L^2 forms

$$L^2 \mathcal{L}_{p,q}^k := \bigoplus_{r+s=k, r < p, s < q} L^2 \Lambda^{r,s} \quad \text{for } k \leq p + q - 2,$$

$$L^2 \mathcal{L}_{p,q}^{k-1} := \bigoplus_{r+s=k, r \geq p, s \geq q} L^2 \Lambda^{r,s} \quad \text{for } k \leq p + q,$$

so that for $a_k \in \{s, w\}$ with $a_k \leq a_{k+1}$ we have closed and densely defined extensions $(\delta_{p,q}^k)_{a_k}$. The full L^2 ABC Hilbert complex is therefore

$$0 \longrightarrow L^2 \mathcal{L}_{p,q}^0 \xrightarrow{(\delta_{p,q}^0)_{a_0}} L^2 \mathcal{L}_{p,q}^1 \longrightarrow \dots \longrightarrow L^2 \mathcal{L}_{p,q}^{2n-1} \xrightarrow{(\delta_{p,q}^{2n-1})_{a_{2n-1}}} L^2 \mathcal{L}_{p,q}^{2n} \longrightarrow 0.$$

Remark 10.1. We chose to define the L^2 ABC Hilbert complex using just the strong and the weak extensions for more clarity. We point out that the same definition could be rephrased using the more general notion of ideal boundary condition introduced in [11, Section 3].

We now list the following natural open questions arising in this context.

10.1. *Possibly incomplete Kähler metrics*

In Section 8 we studied the L^2 Hodge theory for the Bott-Chern and Aeppli Laplacians on complex manifolds endowed with a complete Kähler metric. It would be worthwhile to establish which results continue to be true after dropping the assumption of completeness.

10.2. L^2 Frölicher inequality

Given a compact complex manifold, in [3, Theorems A,B] Angella and Tomassini have proved that

$$h_{\bar{\partial}}^{p,q} + h_{\bar{\partial}}^{p,q} \leq h_A^{p,q} + h_{BC}^{p,q},$$

and, taking into account the Frölicher inequality

$$2b^k \leq \sum_{p+q=k} h_{\bar{\partial}}^{p,q} + h_{\bar{\partial}}^{p,q}$$

of [20, Theorem 2], they also proved that

$$2b^k = \sum_{p+q=k} h_A^{p,q} + h_{BC}^{p,q}$$

if and only if the $\bar{\partial}\bar{\partial}$ -Lemma holds. We are therefore motivated to ask the following question.

Does a Frölicher inequality between L^2 Betti and L^2 Hodge numbers

$$b_{\Gamma}^k(M) \leq \sum_{p+q=k} h_{\bar{\partial},\Gamma}^{p,q}(M)$$

hold?

In case of an affirmative answer it would follow that

$$2b_{\Gamma}^k(M) \leq \sum_{p+q=k} h_{A,\Gamma}^{p,q}(M) + h_{BC,\Gamma}^{p,q}(M)$$

and one could study the limit case given by the equality, like in the case of compact complex manifolds.

10.3. L^2 index Theorem for the ABC complex

The L^2 index theorem of Atiyah [5] states that on a Galois Γ -covering of a compact manifold M , the analytical index of an elliptic operator P between vector bundles coincides with the Γ -index of the lift operator \tilde{P} . When $P = d + d^*$, one finds that the Euler characteristic can be computed via L^2 Betti numbers

$$\chi(M) = \sum_k (-1)^k b_{\Gamma}^k.$$

If M is a complex manifold and $P = \bar{\partial} + \bar{\partial}^*$, a similar result holds (cf. [7]).

In our setting the natural operator to choose seems to be $P = \sum_k \delta_{p,q}^{2k} + \sum_k (\delta_{p,q}^{2k+1})^*$, however this is not elliptic. Therefore one would first need to generalise the L^2 index Theorem of Atiyah in the case of elliptic complexes.

We refer to [40] for a recent result yielding the equality between the analytical and the topological indexes of the ABC complex $\mathcal{L}_{p,q}^{\bullet}$ on compact complex manifolds.

Data availability

No data was used for the research described in the article.

References

- [1] P. Albin, An Introduction to L^2 -Cohomology, Lecture Notes, London Summer School and Workshop: The Sen Conjecture and Beyond, 2017, <https://ckottke.ncf.edu/senworkshop/IntroL2Coho.pdf>.
- [2] A. Andreotti, E. Vesentini, Carleman estimates for the Laplace-Beltrami equation on complex manifolds, Publ. Math. Inst. Hautes Études Sci. 25 (1965) 81–130.
- [3] D. Angella, A. Tomassini, On the $\partial\bar{\partial}$ -lemma and Bott-Chern cohomology, Invent. Math. 192 (2013) 71–81.
- [4] M.F. Atiyah, The index of elliptic operators on compact manifolds, Sémin. N. Bourbaki 8 (253) (1964) 159–169.

- [5] M.F. Atiyah, Elliptic operators, discrete groups and von Neumann algebras, in: Colloque “Analyse et Topologie” en l’Honneur de Henri Cartan, Orsay, 1974, in: Astérisque, vols. 32–33, Soc. Math. France, Paris, 1976, pp. 32–33.
- [6] F. Bei, On the L^2 -Poincaré duality for incomplete Riemannian manifolds: a general construction with applications, *J. Topol. Anal.* 8 (1) (2016) 151–186.
- [7] F. Bei, Von Neumann dimension, Hodge index theorem and geometric applications, *Eur. J. Math.* 5 (4) (2019) 1212–1233.
- [8] F. Bei, Symplectic manifolds, L^p -cohomology and q-parabolicity, *Differ. Geom. Appl.* 64 (2019) 136–157.
- [9] F. Bei, S. Diverio, P. Eyssidieux, S. Trapani, Weakly Kähler hyperbolic manifolds and the Green–Griffiths–Lang conjecture, *J. Reine Angew. Math.* 807 (2024) 257–297.
- [10] B. Bigolin, Gruppi di Aeppli, *Ann. Sc. Norm. Super. Pisa, Cl. Sci.* (3) 23 (1969) 259–287.
- [11] J. Brüning, M. Lesch, Hilbert complexes, *J. Funct. Anal.* 108 (1) (1992) 88–132.
- [12] P.R. Chernoff, Essential self-adjointness of powers of generators of hyperbolic equations, *J. Funct. Anal.* 12 (1973) 401–414.
- [13] P. Deligne, P. Griffiths, J. Morgan, D. Sullivan, Real homotopy theory of Kähler manifolds, *Invent. Math.* 29 (3) (1975) 245–274.
- [14] J.P. Demailly, Estimations L^2 pour l’opérateur $\bar{\partial}$ d’un fibré vectoriel holomorphe semi-positif au-dessus d’une variété Kählérienne complète, *Ann. Sci. Éc. Norm. Supér.* (4) 15 (3) (1982) 457–511.
- [15] J.P. Demailly, *Complex Analytic and Differential Geometry*, Université de Grenoble, Saint-Martin d’Hères, 2012.
- [16] H. Donnelly, C. Fefferman, L^2 -cohomology and index theorem for the Bergman metric, *Ann. Math.* 118 (1983) 593–618.
- [17] B. Eckmann, Introduction to L^2 -methods in topology: reduced L^2 homology, harmonic chains, L^2 -Betti numbers, *Isr. J. Math.* 117 (2000) 183–219.
- [18] W.G. Faris, *Self-Adjoint Operators*, Lecture Notes in Mathematics, vol. 433, Springer-Verlag, Berlin-New York, 1975, vii+115 pp.
- [19] G.B. Folland, J.J. Kohn, The Neumann Problem for the Cauchy-Riemann Complex, *Annals of Mathematics Studies*, vol. 75, Princeton University Press/University of Tokyo Press, Princeton, N.J./Tokyo, 1972, viii+146 pp.
- [20] A. Frölicher, Relations between the cohomology groups of Dolbeault and topological invariants, *Proc. Natl. Acad. Sci. USA* 41 (1955) 641–644.
- [21] M. Gromov, Kähler hyperbolicity and L_2 -Hodge theory, *J. Differ. Geom.* 33 (1991) 263–292.
- [22] B.C. Hall, *Quantum Theory for Mathematicians*, Graduate Texts in Mathematics, vol. 267, Springer, New York, 2013, xvi+554 pp.
- [23] R.K. Hind, A. Tomassini, On L_2 -cohomology of almost Hermitian manifolds, *J. Symplectic Geom.* 17 (6) (2019) 1773–1792.
- [24] L. Hörmander, L^2 estimates and existence theorems for the $\bar{\partial}$ operator, *Acta Math.* 113 (1965) 89–152.
- [25] T. Huang, Q. Tan, L^2 -hard Lefschetz complete symplectic manifolds, *Ann. Mat. Pura Appl.* (4) 200 (2) (2021) 505–520.
- [26] K. Kodaira, D.C. Spencer, On deformations of complex analytic structures, III. Stability theorems for complex structures, *Ann. Math.* 71 (1960) 43–76.
- [27] M. Lesch, Differential operators of Fuchs type, conical singularities, and asymptotic methods, preprint, arXiv:dg-ga/9607005, 1996.
- [28] W. Lück, L^2 -Invariants: Theory and Applications to Geometry and K -Theory, *Ergebnisse der Mathematik und Ihrer Grenzgebiete, 3, Folge (A Series of Modern Surveys in Mathematics)*, vol. 44, Springer-Verlag, Berlin, ISBN 3-540-43566-2, 2002, xvi+595 pp.
- [29] S. Marouani, D. Popovici, Balanced hyperbolic and divisorially hyperbolic compact complex manifolds, *Math. Res. Lett.* 30 (6) (2023) 1813–1855.
- [30] S. Marouani, D. Popovici, Some properties of balanced hyperbolic compact complex manifolds, *Int. J. Math.* 33 (3) (2022) 2250019, 39 pp.
- [31] V. Pati, Elliptic complexes and index theory, Lecture Notes, <https://www.isibang.ac.in/~adean/infsys/database/notes/elliptic.pdf>.
- [32] R. Piovani, $W^{1,2}$ Bott-Chern and Dolbeault decompositions on Kähler manifolds, *J. Geom. Anal.* 33 (281) (2023).
- [33] R. Piovani, A. Tomassini, Bott-Chern harmonic forms on Stein manifolds, *Proc. Am. Math. Soc.* 147 (2019) 1551–1564.

- [34] R. Piovani, A. Tomassini, Bott-Chern harmonic forms on complete Hermitian manifolds, *Int. J. Math.* 30 (5) (2019).
- [35] R. Piovani, A. Tomassini, Aeppli cohomology and Gauduchon metrics, *Complex Anal. Oper. Theory* 14 (1) (2020) 22, 15 pp.
- [36] M. Reed, B. Simon, *Methods of Modern Mathematical Physics. I. Functional Analysis*, second edition, Academic Press, New York, 1972.
- [37] M. Reed, B. Simon, *Methods of Modern Mathematical Physics. II. Fourier Analysis, Self-Adjointness*, Academic Press, New York-London, 1975.
- [38] J. Ruppenthal, L^2 -theory for the $\bar{\partial}$ -operator on complex spaces with isolated singularities, *Ann. Fac. Sci. Toulouse Math.* (6) 28 (2) (2019) 225–258.
- [39] M. Schweitzer, Autour de la cohomologie de Bott-Chern, preprint, arXiv:0709.3528v1, 2007.
- [40] J. Stelzig, Some remarks on the Schweitzer complex, preprint, arXiv:2204.06027, 2022.
- [41] R.S. Strichartz, Analysis of the Laplacian on the complete Riemannian manifold, *J. Funct. Anal.* 52 (1) (1983) 48–79.
- [42] Q. Tan, H. Wang, J. Zhou, Symplectic parabolicity and L^2 symplectic harmonic forms, *Q. J. Math.* 70 (1) (2019) 147–169.
- [43] J. Varouchas, Propriétés cohomologiques d’une classe de variétés analytiques complexes compactes, in: P. Lelong, P. Dolbeault, H. Skoda (Eds.), *Séminaire d’analyse, années 1983/1984*, in: *Lecture Notes in Math.*, vol. 1198, Springer, Berlin, 1986, pp. 233–243.